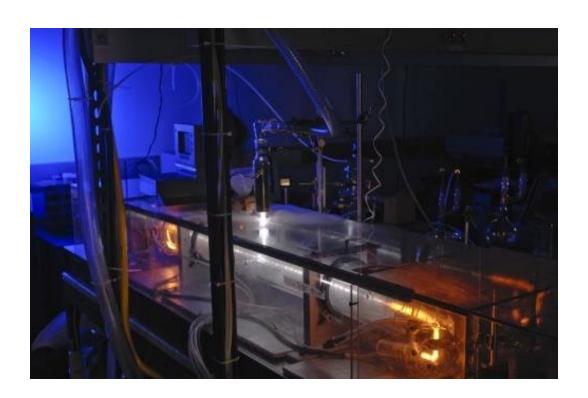


### Seminar at Department of Aerospace Engineering, University of Michigan January 30, 2020

## **Understanding Reactivity of Nonequilibrium Molecular Plasmas for Propulsion and Power Applications**

#### Igor Adamovich

Nonequilibrium Thermodynamics Laboratory (NETL)
Department of Mechanical and Aerospace Engineering





#### Nonequilibrium Thermodynamics Laboratory: Research Fields

- Kinetics of high-speed nonequilibrium flows and low-temperature plasmas
- Energy transfer in molecular collisions, nonequilibrium chemical reactions, emission and absorption of radiation
- Laser diagnostics of plasmas and reacting flows
- Development of new molecular gas lasers
- Applications for high-speed aerodynamics, propulsion and combustion, hypersonic flows, plasma chemical reactors, and biology



#### Nonequilibrium Reacting Plasmas: Relevance for Aerospace Applications

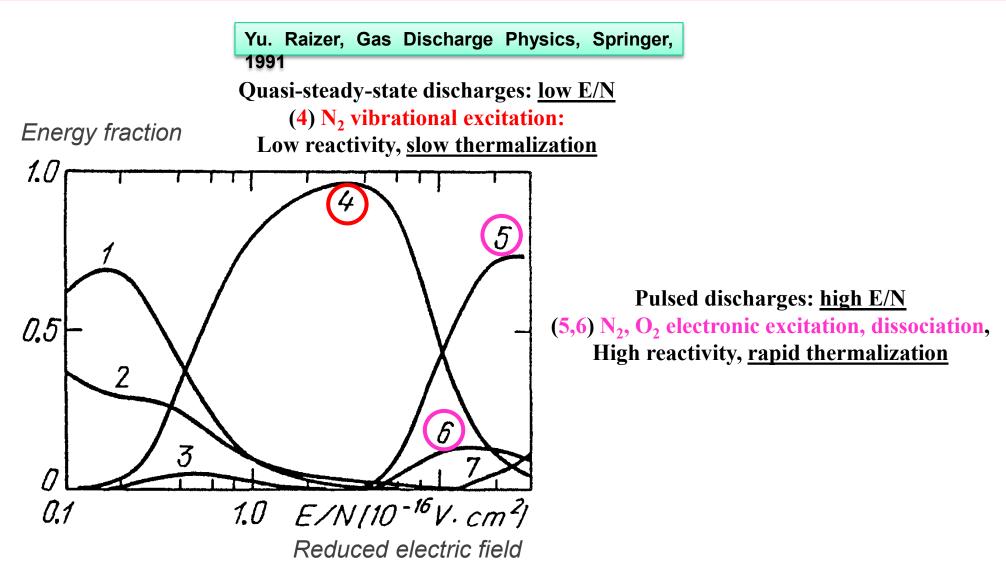
- High-speed flow (formation of coherent structures) is affected by localized heating in pulsed plasmas: <u>plasma flow control</u>
- Excited species and radicals enable low-temperature reaction pathways in reacting flows: plasma-assisted combustion
- Excited species and radicals control UV / visible emission from nonequlibrium flows: hypersonic flight, atmospheric reentry
- Energy partition in the plasma (vibrational and electronic excitation, dissociation, radical generation), and temperature rise are controlled by the electric field
- Quantitative insight requires non-intrusive diagnostics of electric field, excited species, and radicals:
  - In laboratory environments
  - At pulsed high-enthalpy flow facilities



# Electric Field in Pulsed Plasmas: Understanding High-Speed Plasma Flow Control



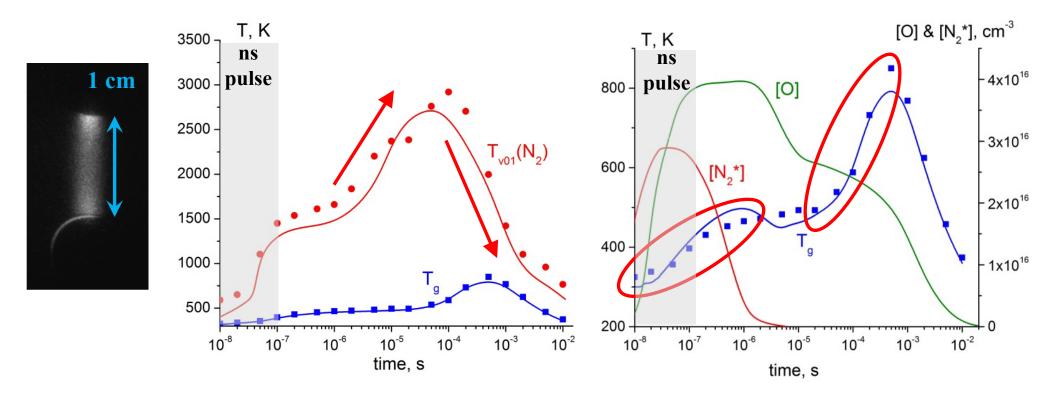
#### **Energy Partition in Air Plasma vs. Electric Field**



• E/N controls species generated in the plasma and the rate of temperature rise



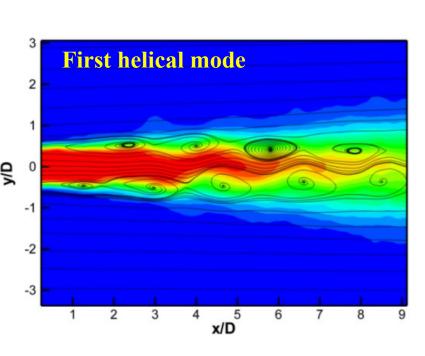
### Effect of Excited Species on Temperature Rise: Ns Pulse Discharge in Air, 100 Torr

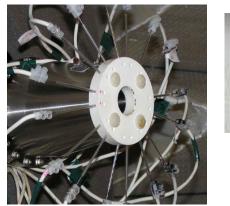


- $T_v$  rise: V-V exchange,  $N_2(v) + N_2(v=0) \rightarrow N_2(v-1) + N_2(v=1)$
- $T_v$  decay: V-T relaxation,  $N_2(v) + O \rightarrow N_2(v-1) + O$
- "Rapid" heating: quenching of  $N_2$  electronic states,  $N_2$ \* +  $O_2 \rightarrow N_2$  + O + O
  - Compression wave formation
- "Slow" heating: V-T relaxation,  $N_2(v) + O \rightarrow N_2(v-1) + O$ 
  - Formation of coherent structures in the flow

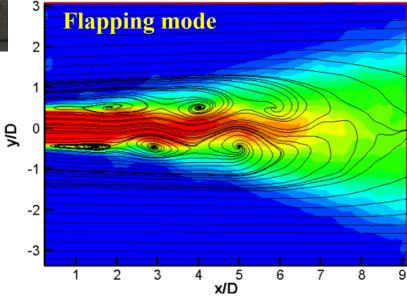


## Localized Arc Filament Plasma Flow Actuators M=0.9 Circular Jet







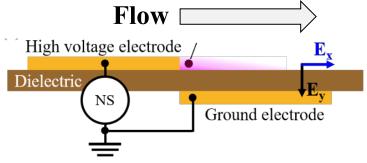


- High amplitude perturbations (localized heating in arc filaments)
- Every discharge pulse results in a vortex formation
- Flow responds to forcing near instability frequency
- Mixing enhancement, jet noise reduction

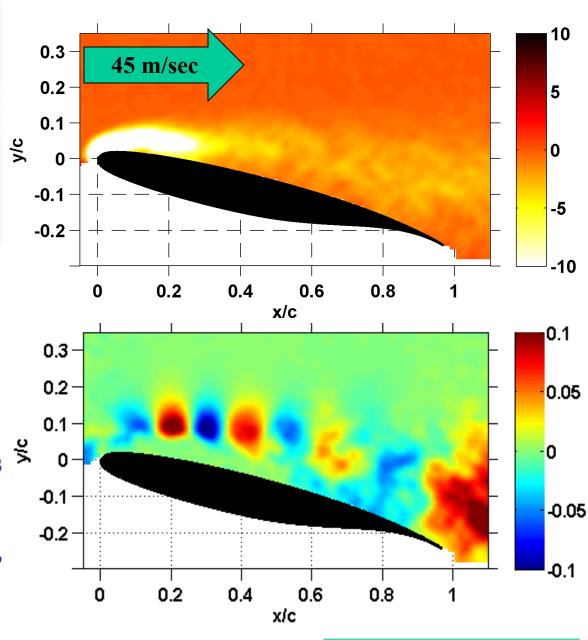


## Ns Pulse Surface Plasma Flow Actuators: Boundary Layer Reattachment (M=0.2-0.3)





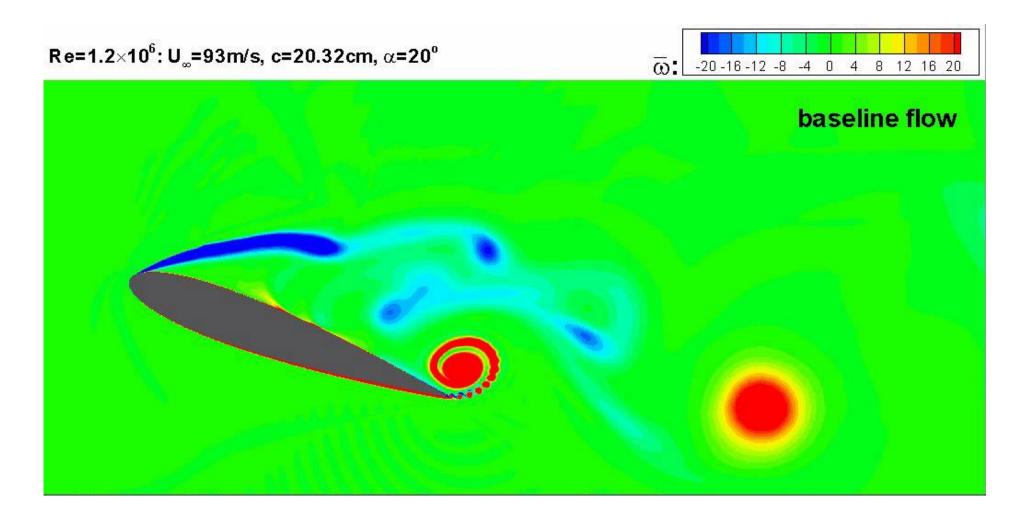
- Every ns discharge pulse produces sa spanwise "vortex tube"
- Enhanced mixing with free stream, boundary layer reattachment
- Effect detected up to u=96 m/s



Little et al, AIAA J., 2012



### **Surface Plasma / CFD Modeling Predictions Baseline and forced flows**



- Vortex formation controlled by localized heating in plasma layer
- Need to know time-resolved electric field, E(t), to predict flow field accurately

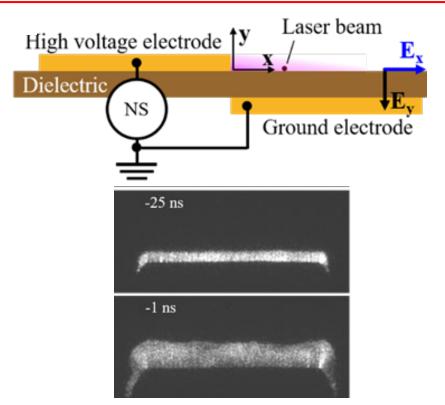


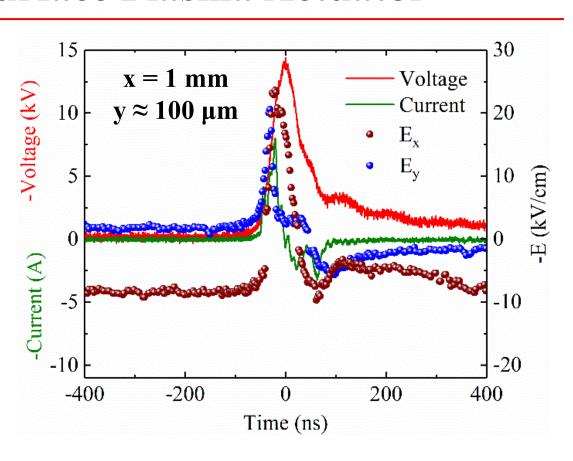
## Electric Field Induced Second Harmonic Generation (E-FISH)

- Electric field in the plasma, E, induces a dipole moment in molecules or atoms
- Pulsed laser beam (fs, ps, or even ns) passes through the plasma
- Laser field,  $I_{\omega}$ , generates coherent oscillating polarization of the dipoles
- Oscillating dipoles launch a coherent  $2^{\rm nd}$  harmonic signal beam,  $I_{2\omega} \sim \left[\chi^3 I_{\omega} EL\right]^2$   $(\chi^{(3)}$  nonlinear susceptibility, L interaction length)
- Straightforward signal isolation and detection, sub-ns time resolution
- Signal polarization same as field direction:  $E_x$ ,  $E_y$  can be measured separately
- Simplest experiment ever, great for a student laser diagnostics lab
- Goal: determine effect of discharge pulse waveform on the flow field



## Electric Field In Ns Pulse Surface Plasma Actuator

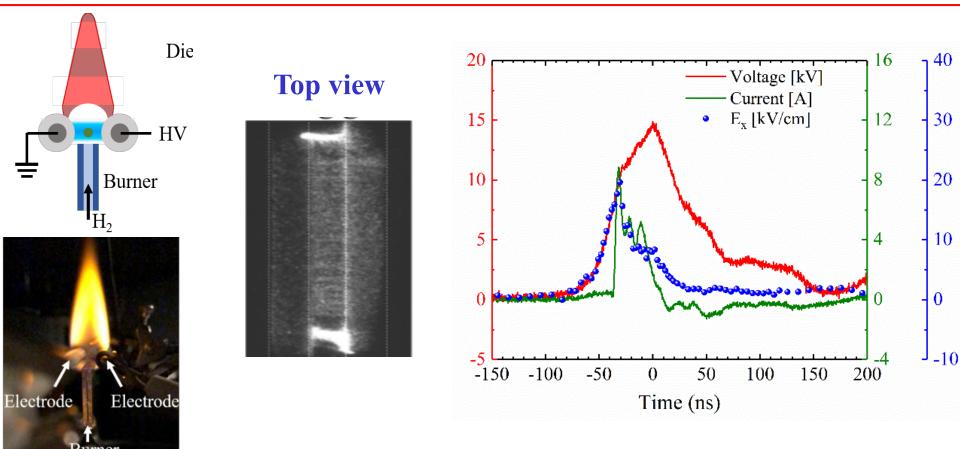




- 150 ps Nd:YAG pulse, 2 mJ at 1064 nm
- Field offset: surface charge accumulation from previous discharge pulse
- Field rises with applied voltage, drops after breakdown (plasma self-shielding)
- E(t)/N controls temperature rise, vortex formation in the flow



## Electric Field In H<sub>2</sub> Diffusion Flame with Ns Pulse Discharge



- Electric field measured in  $H_2$  plasma at T=370 K
- Energy coupled at E = 9-19 kV/cm (E/N = 50-100 Td): efficient H atom generation
- E/N(t) is critical for quantifying effect of plasma on combustion

Simeni Simeni et al., Comb. Flame, 2018

Simeni Simeni et al., Comb. Flame, 2019



### "Dark" Metastable Species in Hypersonic Flows: Understanding UV Radiation Mechanism

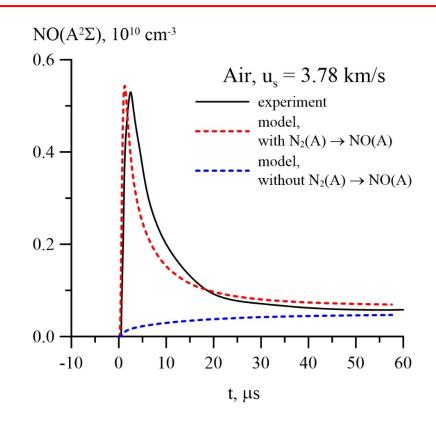


#### **UV Emission from Strong Shock Waves in Air**



• Shock tube data, modeling (CUBRC, 1991): NO UV radiation controlled by energy transfer from  $N_2(A^3\Sigma)$ 

$$N_2 + N \rightarrow N_2(A^3\Sigma) + N$$
 $N_2(A^3\Sigma) + NO \rightarrow N_2 + NO(A^2\Sigma)$ 
 $NO(A^2\Sigma) \rightarrow NO + hv \quad (\gamma \ bands)$ 



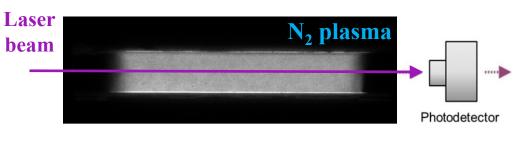
NO UV emission, Mach 11 normal shock in air

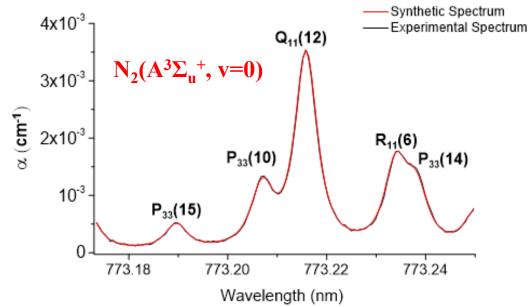
- No one ever detected  $N_2(A^3\Sigma)$  in a hypersonic flow it is a "dark" metastable state
- Goal: diagnostics of "dark" states such as  $N_2(A^3\Sigma)$  and  $O_2(a^1\Delta)$  at high-enthalpy test facilities, isolate UV radiation and  $O_2$  dissociation mechanisms



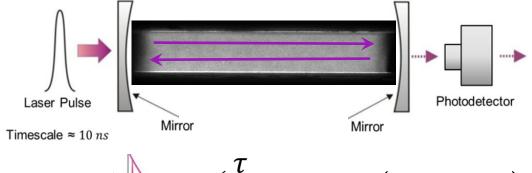
#### **How to Measure Species** That Do Not Want to Be Detected

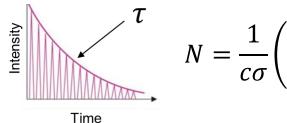
#### **Tunable Diode Laser Absorption Spectroscopy (TDLAS) – single pass**



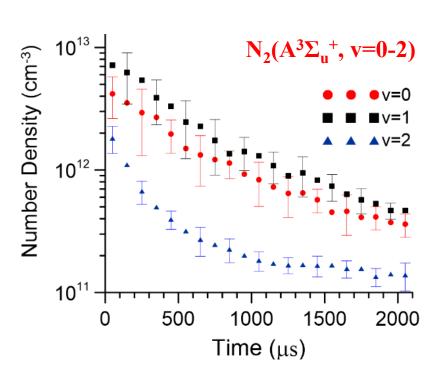


#### **Cavity Ring Down Spectroscopy (CRDS)** ~ 10,000 passes





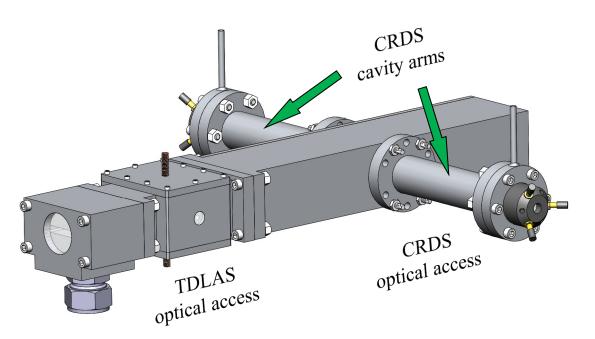
$$N = \frac{1}{c\sigma} \left( \frac{1}{\tau} - \frac{1}{\tau_{empty}} \right)$$



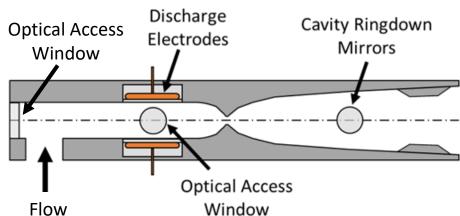
Jans et al, J. Molecular Spectroscopy,



#### Nonequilibrim Flow Mach 5 Wind Tunnel



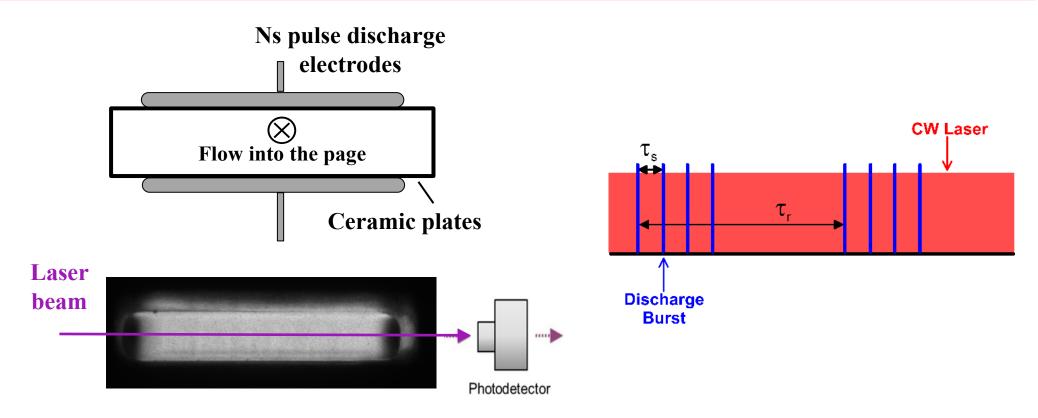




- Blowdown wind tunnel,  $P_0 = 0.3 1.0$  atm, Mach 4 5, run rime 10 s
- $N_2(A)$  generated by discharge in plenum
- Measurements by TDLAS (in plenum) and CRDS (in supersonic flow)



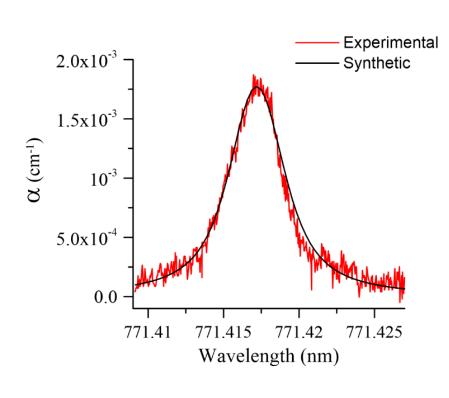
#### **TDLAS Measurements in Plenum**

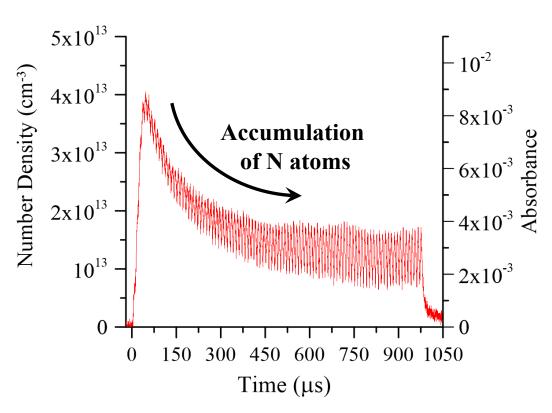


- $N_2(A)$  generated in a repetitive ns pulse discharge
- Diffuse, uniform plasma fills entire flow cross section
- Laser is scanned across absorption line or "parked" at the line center
- Same approach can be used in a shock tube (O atom measurements at Stanford, laser scan takes 1-2  $\mu s$ )

#### **TDLAS** Measurements in Plenum

• Nitrogen,  $P_0 = 250$  Torr,  $T_0 = 320$  K, discharge pulse rep rate v=100 kHz, 100 pulse burst



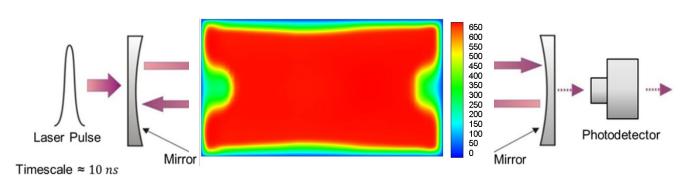


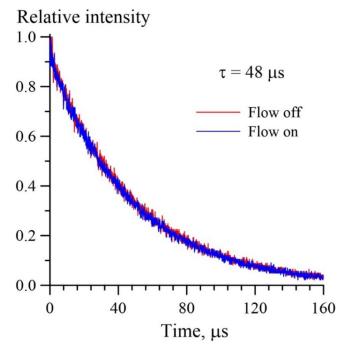
- $N_2(A)$  generated by electron impact:  $N_2 + e \rightarrow N_2(A^3\Sigma, B^3\Sigma, C^3\Pi) + e \rightarrow N_2(A^3\Sigma) + e$
- Quenching dominated by N atoms:  $N_2(A^3\Sigma) + N \rightarrow N_2 + N$
- Behind the shock, excitation of  $N_2$  by N atoms is the dominant  $N_2(A)$  generation process



#### CRDS Measurements in a Supersonic Flow: Can It Be Done?

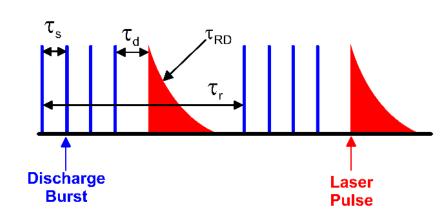
#### Velocity in a Mach 4 test section, CFD predictions (flow into the page)





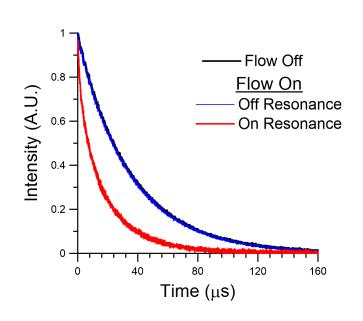
Would side wall boundary layers cause beam steering?

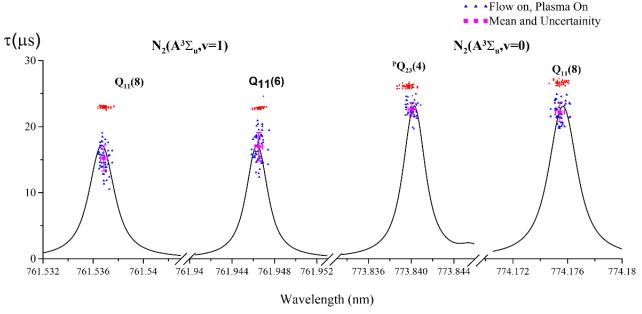
- Nitrogen,  $P_0 = 227$  torr, M = 4.2
- Robust optical setup: no effect of building or wind tunnel vibrations
- No effect of flow on empty cavity ring down time
- Next, turn on discharge bursts in plenum:



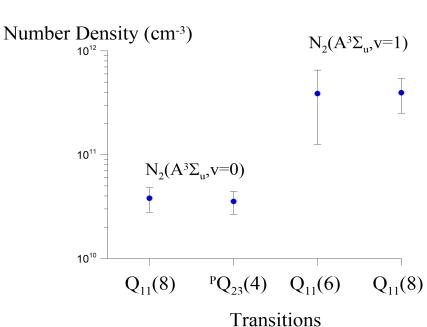


#### **CRDS** Measurements in a Mach 4 Flow





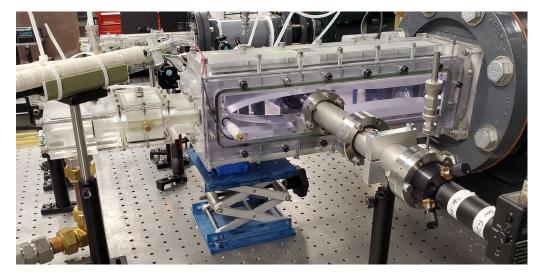
- $P_0 = 227 \text{ torr}, P=2.0 \text{ Torr} (M = 3.7)$
- v=100 kHz ns pulse discharge in plenum
- Single-shot CRDS traces, 50 laser shots
- Ring down time with plasma on is shorter
- $[N_2(A^3\Sigma_u^+,v=0,1)]$  are measured
- Detection limit  $\sim 10^{10}$  cm<sup>-3</sup>

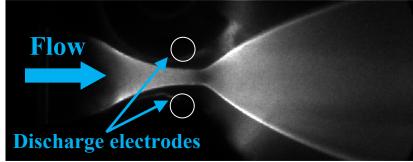


Synthetic SpectrumFlow on, Plasma Off

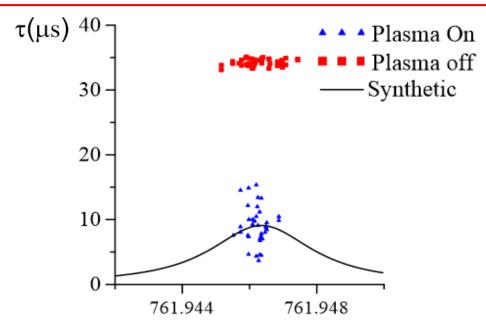


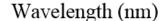
#### **CRDS** Measurements in a Mach 5 Flow

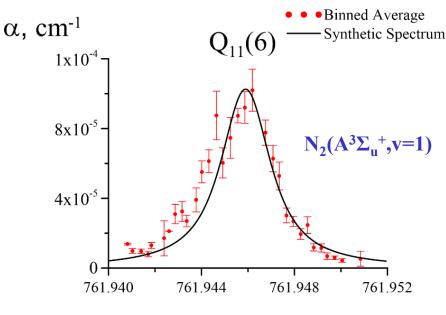




- $N_2$ ,  $P_0 = 250$  Torr, P=1.1 Torr (M = 4.3)
- v=100 kHz discharge at the nozzle throat
- 50 single-shot CRDS traces
- $[N_2(A^3\Sigma_u^+, v=0)] = 2 \cdot 10^{11} \text{ cm}^{-3}, T = 70 \text{ K}$







Wavelength (nm)



#### **Hybrid Plasmas:**

**Understanding Plasma Chemistry** 

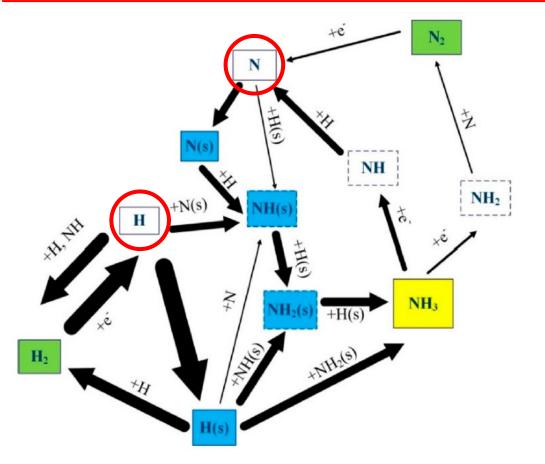
and Plasma Catalysis Mechanisms

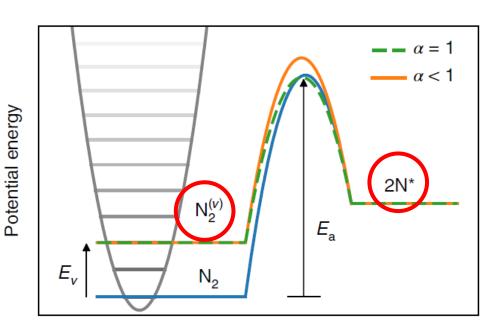
#### "Hybrid" Plasmas

- Separate waveforms for ionization and main energy loading
  - Ionization and energy addition are uncoupled, stable plasmas at high pressures
  - Previously used for efficient molecular lasers (CO<sub>2</sub>, CO, e-COIL)
- High peak E/N ns pulses:
  - Sustain ionization
  - Generate electronically excited species, e.g.  $N_2(A^3\Sigma_u^+)$ , and atoms (N, H, O)
- Main energy loading waveform (DC or RF):
  - Generate vibrationally excited molecules, e.g.  $N_2(v)$ ,  $H_2(v)$ , CO(v),  $CO_2(v_1, v_2, v_3)$
- Goal: isolate the effect of atoms / radicals and vibrationally excited species on plasma chemistry and plasma catalysis
- Goal: develop efficient plasma chemical syntheses at atmospheric pressure



## Plasma Catalytic Synthesis of NH<sub>3</sub>: Reaction Pathways





Reaction coordinate

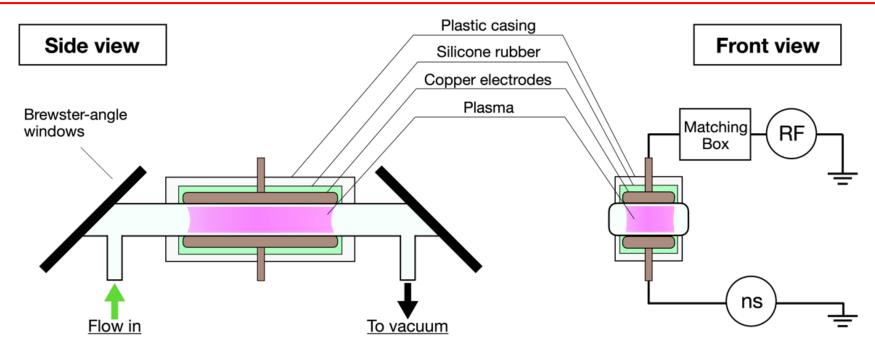
- Process dominated by surface reactions of N and H (generated by electron impact)
- Vibrational excitation of  $N_2$  reduces barrier for surface dissociation reaction (?)
- J. Shah et al, ACS Appl. Energy Mater. 2018

P. Mehta et al., Nature Catalysis 2018

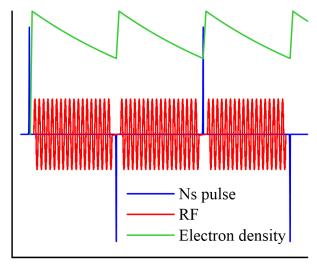
N / H or  $N_2(v)$ ? It is difficult to generate them independently



#### Ns Pulse / RF Hybrid Plasma



Voltage, Electron Density

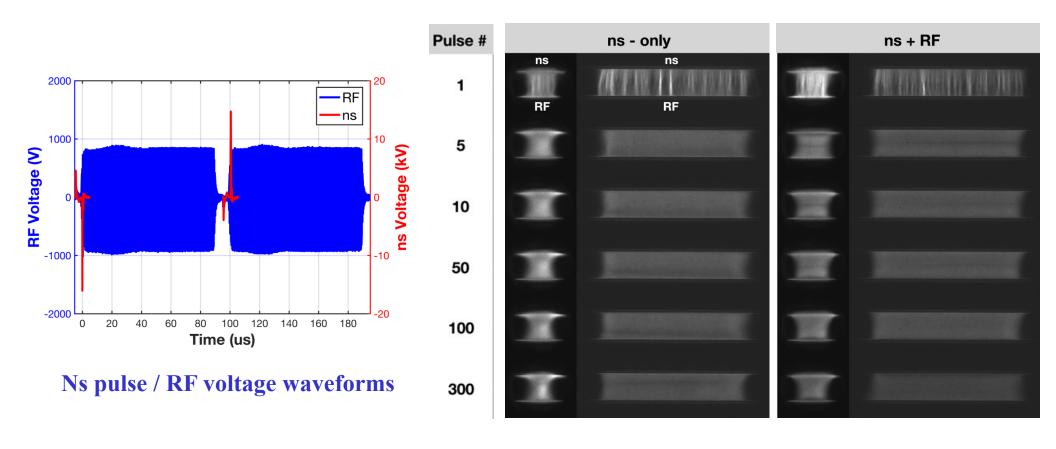


Time

- Single pair of electrodes, external to the reactor
- Ns pulses and RF bursts are separated in time
- RF voltage couples energy to electrons generated by the pulses



#### Discharge Waveforms and Plasma Images

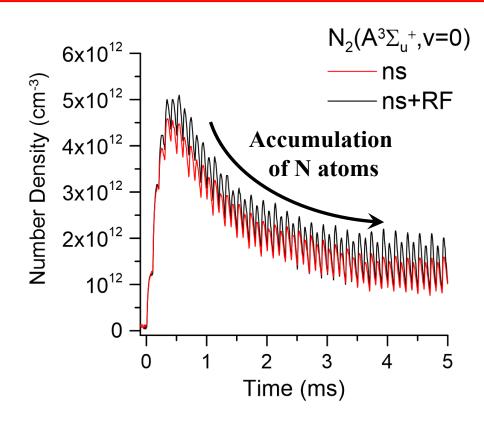


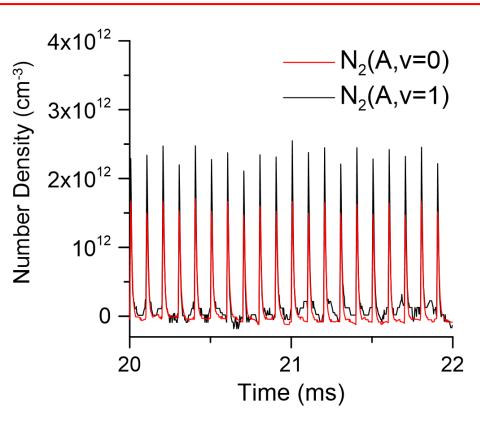
- 100 Torr  $N_2$ , v=10 kHz, diffuse plasma during entire burst
- RF-induced electron oscillations improve plasma uniformity
- Similar behavior up to P=1 atm and v=100 kHz



#### Ns Discharge:

### Generation of $N_2(A^3\Sigma_u^+)$ , N, and H atoms





• [N<sub>2</sub>(A,v=0)] in 100 Torr of nitrogen, with and without RF

- $[N_2(A,v=0,1)]$  in 100 Torr of 1%  $H_2$ - $N_2$ , without RF
- Almost no effect of RF voltage on [N<sub>2</sub>(A)]
- Rapid decay between the pulses

Quenching by N and H atoms,

Quenching by N atoms,

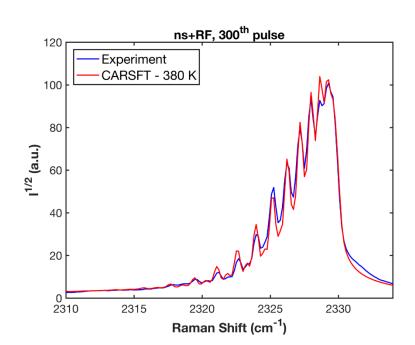
$$N_2(A) + H \rightarrow N_2(X) + H$$

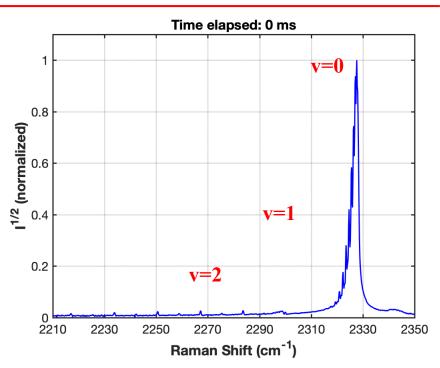
$$N_2(A) + N \rightarrow N_2(X) + N$$

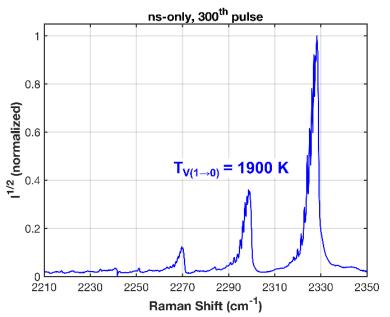


### Ns / RF Discharge: N<sub>2</sub> Vibrational Excitation

- Nitrogen, P=100 Torr, pulse rep rate 10 kHz
- Ns / RF bursts (30 ms on, 70 ms off)
- Animation N<sub>2</sub> CARS spectra
- Strong nonequilibrium: T<sub>v</sub>=1900 K, T=380 K
- Ns and RF: <u>selective generation</u> of excited species









### Summary

- Electric field in pulsed plasmas: flow control *via* quenching of excited species, localized heating, coherent structure formation
- Field measured in plasma flow actuators,  $H_2$ -air plasmas (ps EFISH): quantification of plasma flow control, plasma-assisted combustion
- $N_2(A^3\Sigma)$ , "dark" species controlling NO UV emission behind shock waves: measured in Mach 4-5 nonequlibrium flows (TDLAS and CRDS)
- TDLAS and CRDS: potential for  $N_2(A^3\Sigma)$  and  $O_2(a^1\Delta)$  measurements in shock tubes and shock tunnels, to understand NO UV radiation and  $O_2$  dissociation
- Hybrid plasmas: selective generation of excited species and radicals, potential for isolation of plasma chemistry and catalysis mechanisms
- $N_2(A)$  and  $N_2(v)$ : generated selectively and measured in hybrid ns / RF plasmas (TDLAS, CARS). N, H, product species measurements are underway



### Acknowledgments

#### **Students and Colleagues:**

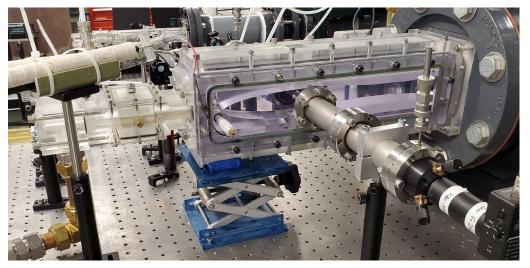
• Ben Goldberg (OSU / Princeton U. / Sandia Livermore), Marien Simeni Simeni (OSU / U. Minnesota), Tang Yong (OSU / Tsinghua U.), Elijah Jans (OSU), Kraig Frederickson (OSU / US Navy), Ilya Gulko (OSU), Bill Rich (OSU), Walter Lempert (OSU), Mo Samimy (OSU), Terry Miller (OSU Chemistry)

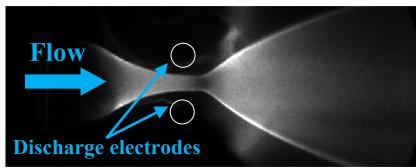
#### **Sponsors:**

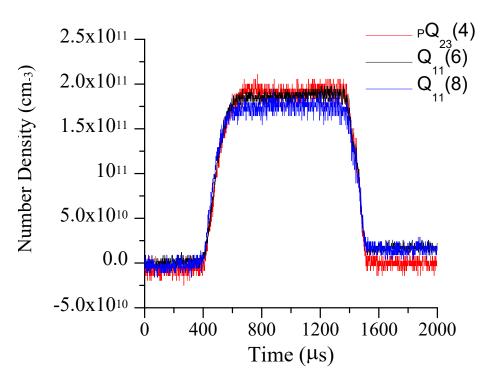
- AFOSR "Energy Transfer Processes in Nonequilibrium Hypersonic Flows"
- US DOE Plasma Science Center "Predictive Control of Plasma Kinetics: Multi-Phase and Bounded Systems"
- NSF "Fundamental Studies of Accelerated Low Temperature Combustion Kinetics by Nonequilibrium Plasmas"
- US DOE PSAAP-2 Center "Exascale Simulation of Plasma-Coupled Combustion"
- US DOE Collaborative Research Center for Studies of Plasma-Assisted Combustion and Plasma Catalysis
- US DOE Center for Low Temperature Plasma Interactions with Complex Interfaces



#### This Just In: $N_2(A^3\Sigma_{11}^+)$ TDLAS in a Mach 5 Flow







- $N_2$ ,  $P_0 = 250$  Torr, P=1.1 Torr (M = 4.3)
- v=100 kHz discharge at the nozzle throat, discharge burst 1 ms long
- $[N_2(A^3\Sigma_u^+, v=0)] = 1.7 \pm 0.1 \cdot 10^{11} \text{ cm}^{-3}, T = 66 \pm 8 \text{ K}$
- Results consistent with CRDS measurements at the same conditions