

Energy conversion in transient molecular plasmas:

Implications for plasma flow control and plasma assisted combustion

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Research backgrounds of students and post-docs: Mechanical Engineering, Aerospace Engineering, Electrical Engineering, Chemical Physics, Physical Chemistry, Physics

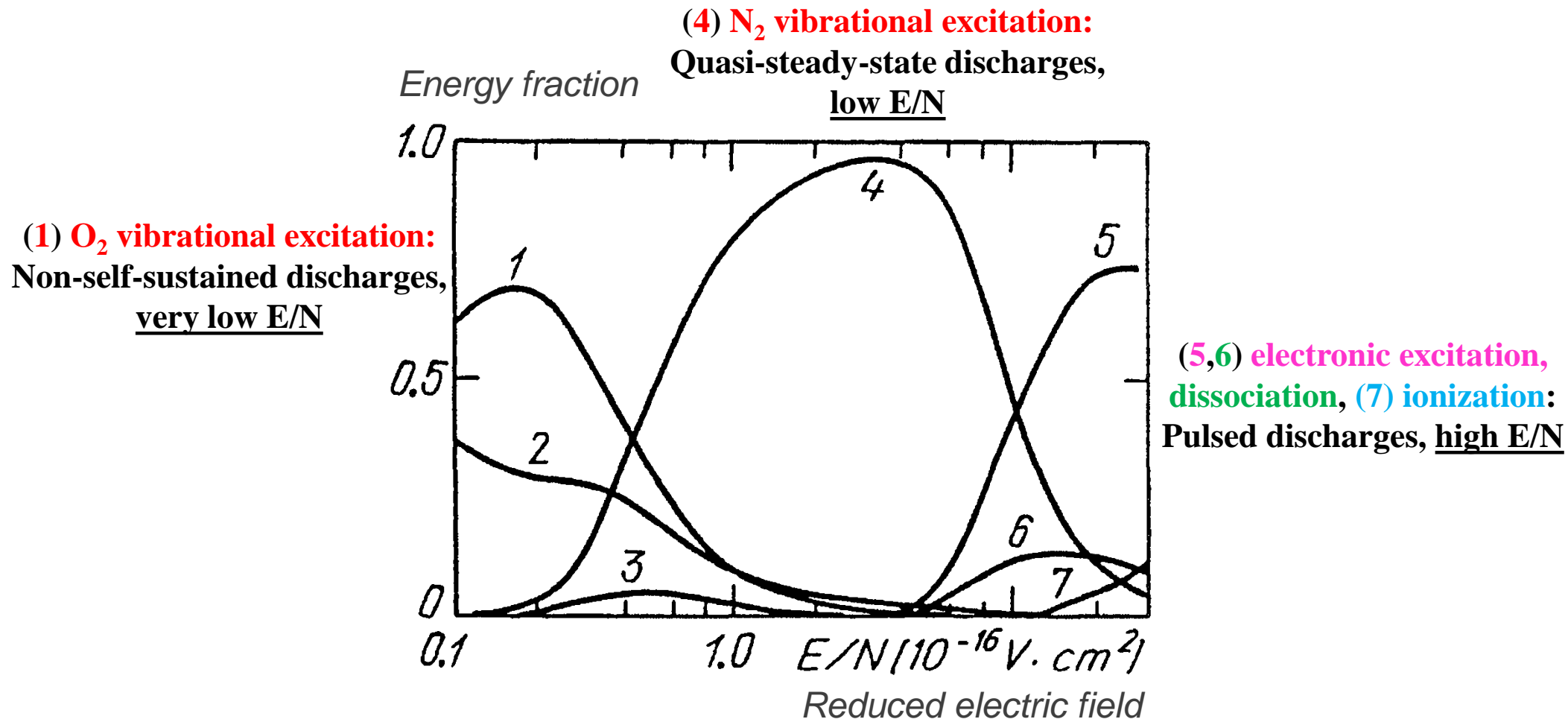
Outline

- I. Motivation: critical importance of energy transfer processes in nonequilibrium, high-pressure molecular plasmas**
- II. Electric field: discharge energy loading and partition**
- III. Electron density and electron temperature: discharge energy loading and partition**
- IV. Dynamics of temperature rise: “rapid” heating and “slow” heating**
- V. Air and fuel-air plasma chemistry: kinetics of plasma assisted combustion**
- VI. Air plasma kinetics and plasma flow control**
- VII. Summary and future outlook**

I. Motivation

Energy Partition in Air Plasma vs. Electric Field

Yu. Raizer, Gas Discharge Physics, Springer, 1991



- Reduced electric field, E/N , controls input energy partition in the discharge
- Rates of electron impact processes: strongly (exponentially) dependent on E/N

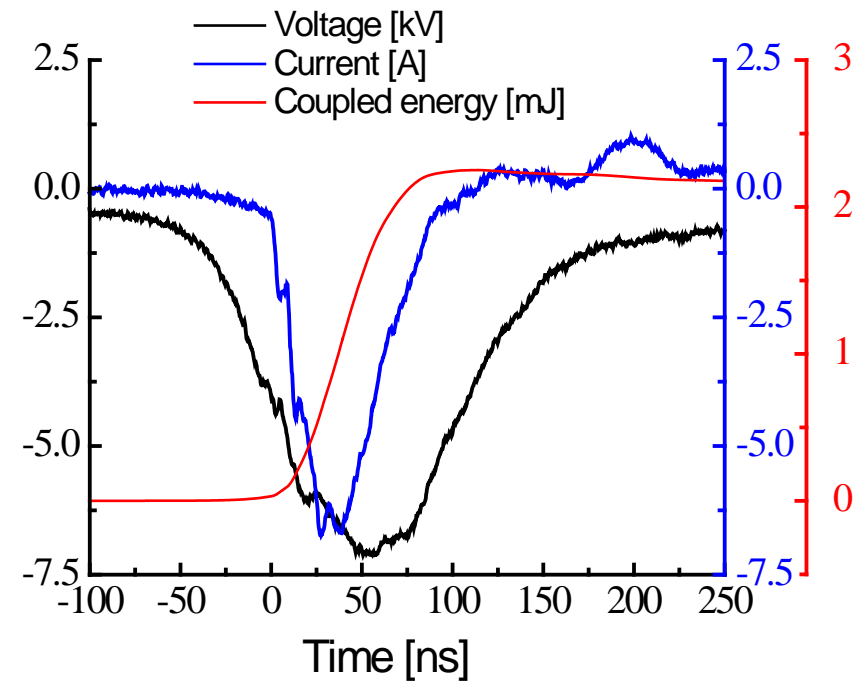
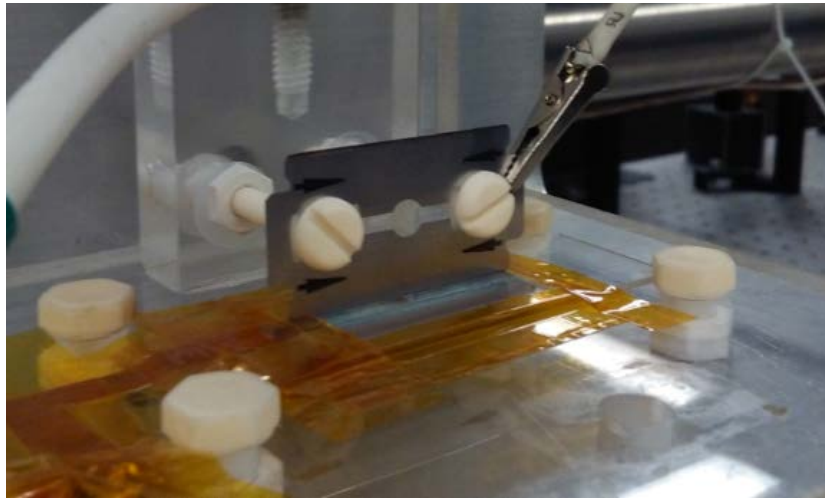
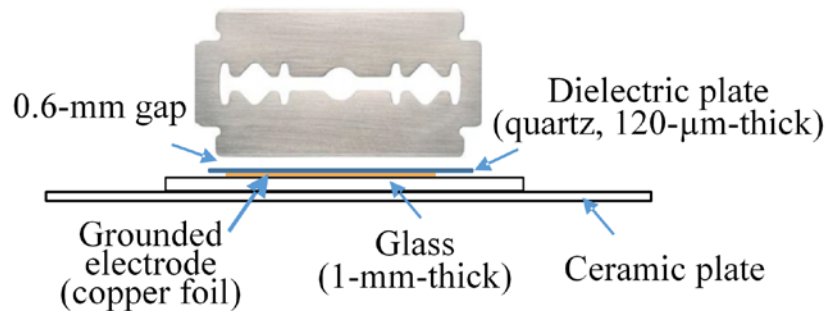
Energy conversion in molecular plasmas: here is what we know

- Energy is coupled to electrons and ions by applied electric field
- Electric field in the plasma: controlled by electron and ion transport, and by surface charge accumulation
- Energy partition (**vibrational** and **electronic** excitation, **dissociation**, **ionization**): controlled by electron density and electric field (or electron temperature)
- Temperature rise in discharge afterglow: controlled by quenching of excited electronic states, vibrational relaxation
- Plasma chemical reactions, rates of radical species generation: controlled by populations of excited electronic states, e.g. N_2^* , excited vibrational states, e.g. $N_2(v)$
- Time-resolved measurements of \vec{E} , n_e , T_e , N_2^* , $N_2(v)$, and radical species (O , H , OH , NO , CH , HO_2 , CH_2O): stable, reproducible, high-pressure ns pulse discharges
- Objective: quantitative insight into energy conversion mechanisms critical for plasma-assisted combustion and plasma flow control

II. Electric field in transient plasmas: insight into discharge energy loading and partition

Diagnostics: CARS-like 4-wave mixing

2-D Ns Pulse Discharge in Atmospheric Air

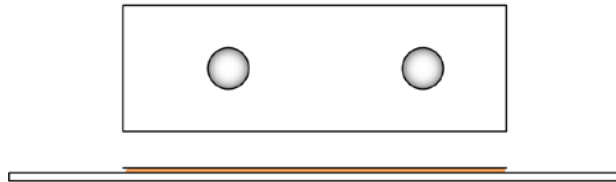


- Discharge sustained between a high-voltage electrode (razor blade) and grounded copper foil, covered with quartz plate 120 μm thick
- Discharge gap 600 μm
- Simple two-dimensional geometry, diffuse plasma

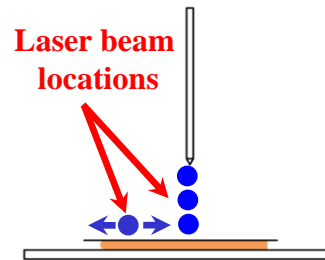
- Peak voltage 7.5 kV, peak current 7 A, coupled energy 2 mJ
- Two current peaks of opposite polarity: “forward” and “reverse” breakdowns
- Time-resolved electric field measured at several locations in the plane of symmetry
- Electric field distribution along the surface is also measured

“Curtain Plasma” Images, Negative Polarity Pulse

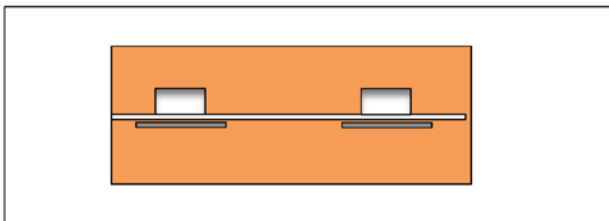
Front View



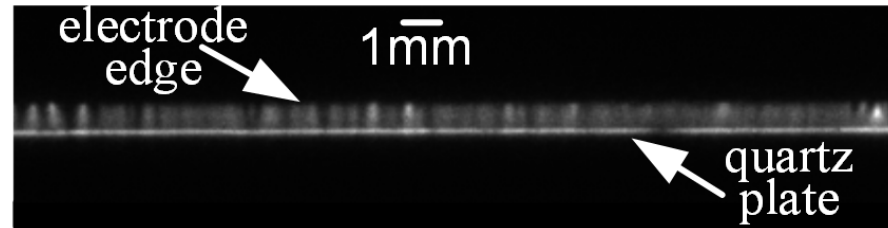
Side View



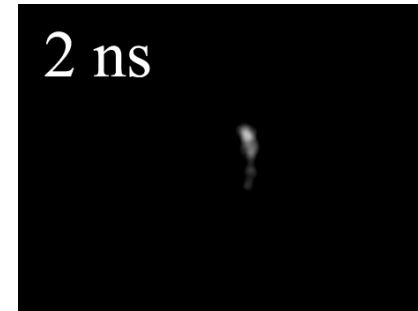
Top View



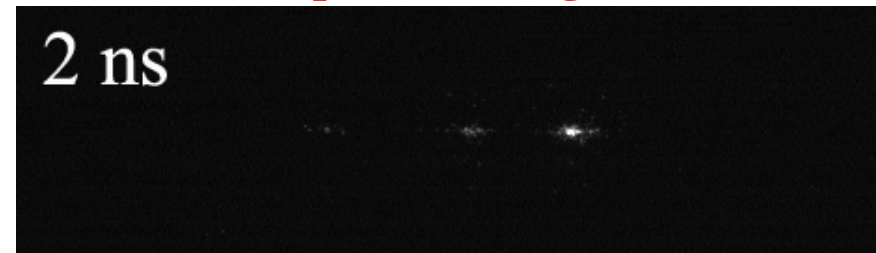
Front view, 100 ns gate



Side view, 2 ns gate

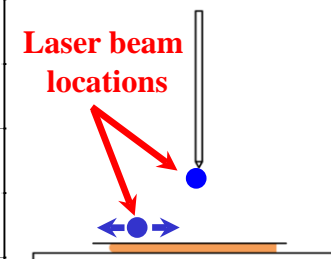
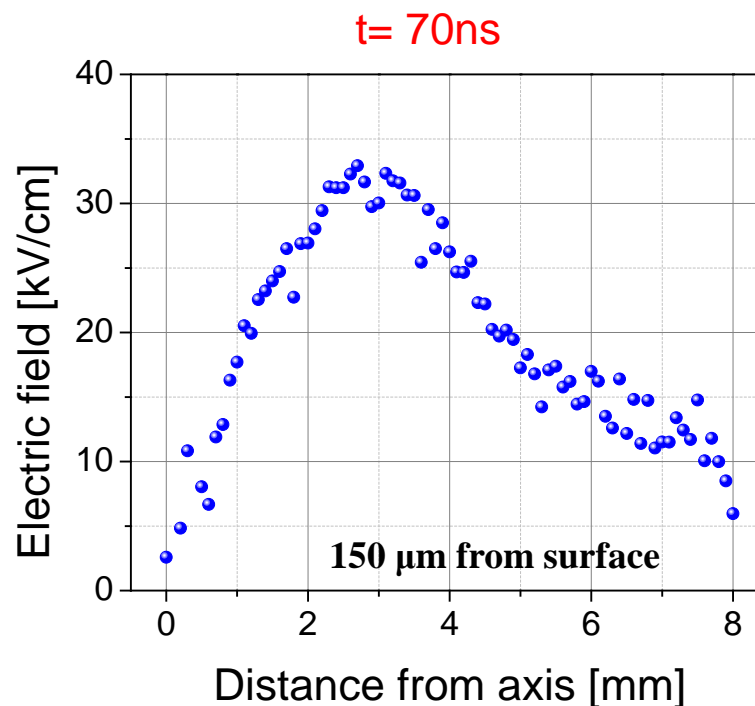
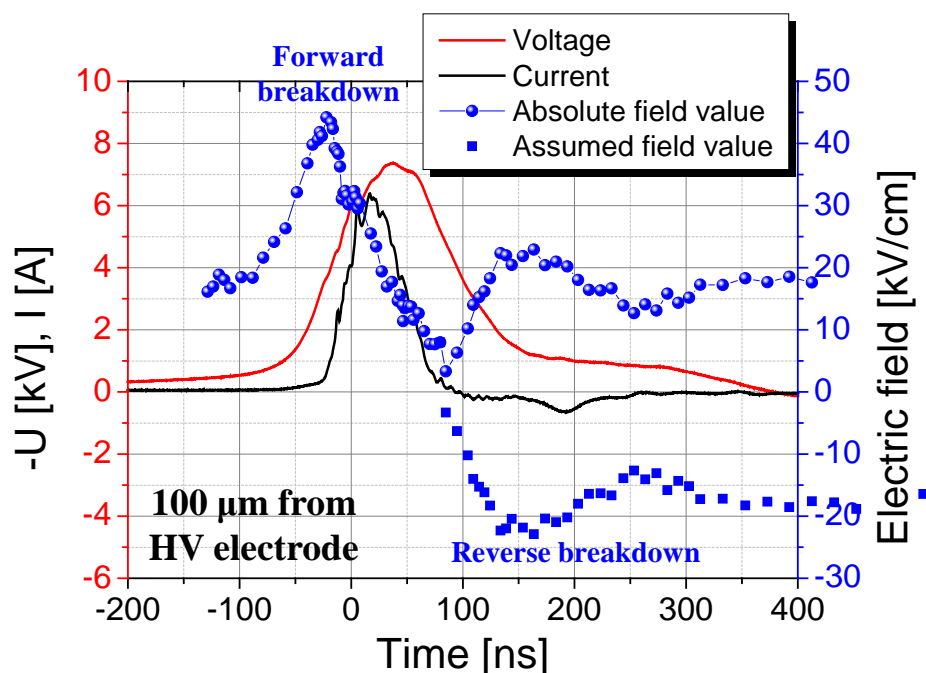


Top view, 2 ns gate



- Front view: near-diffuse plasma “curtain”
- Diffuse surface ionization wave detected, straight ionization front
- Wave speed ~ 0.03 mm/ns
- Surface plasma layer thickness ~ 150 μ m

Time-resolved and Spatially-Resolved Electric Field Measurements

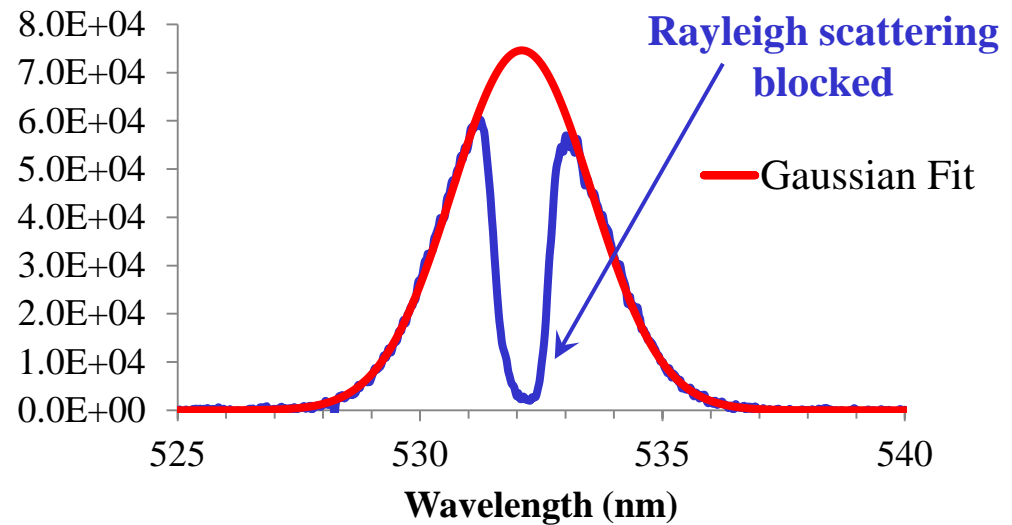
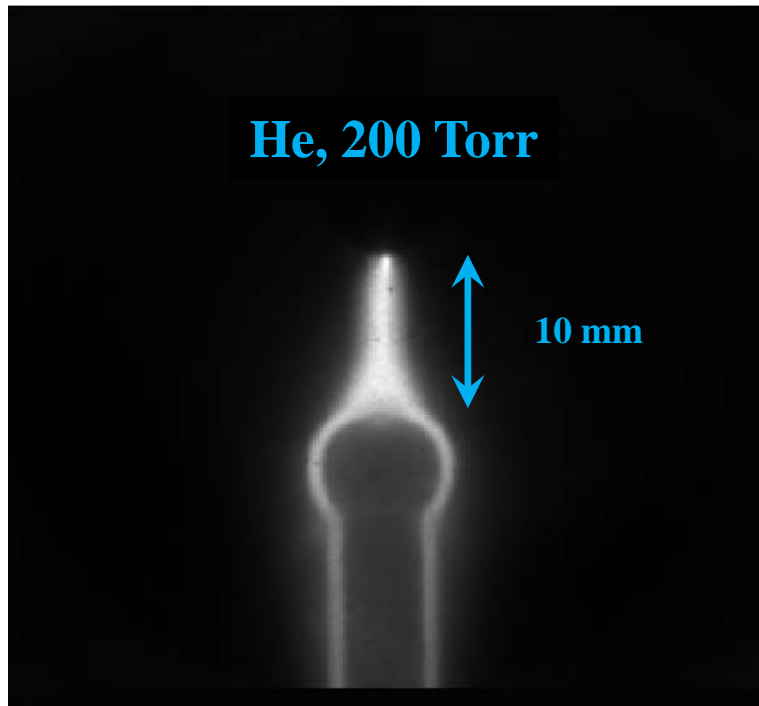


- Initial field offset: charge accumulation on dielectric surface from previous pulse
- Field follows applied voltage rise, increases until “forward breakdown”
- After breakdown, field reduced due to charge accumulation on dielectric surface
- Field is reversed after applied voltage starts decreasing
- After discharge pulse, field decays over several μ s: surface charge neutralization by charges from plasma
- Field distribution measured at the moment when field reversal occurs near HV electrode ($t=70$ ns)
- “Snapshot” electric field distribution across the surface ionization wave front

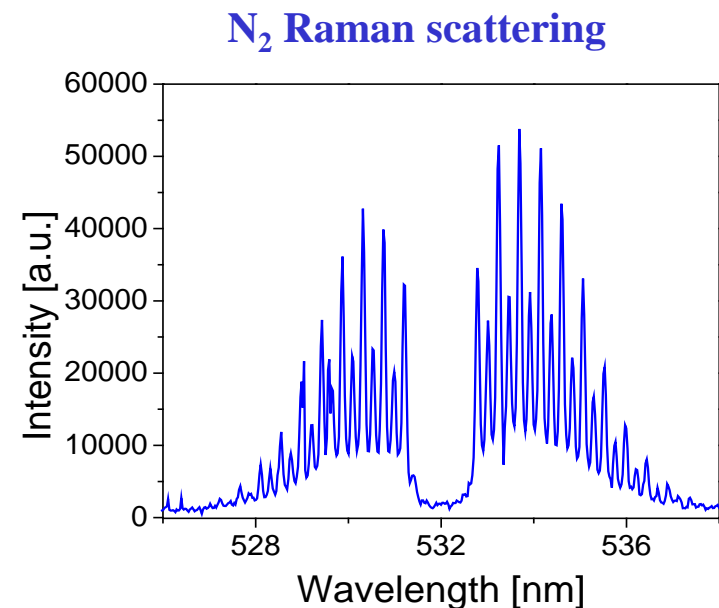
**III. Electron density and electron temperature
in transient plasmas:
insight into discharge energy loading and partition**

Diagnostics: Thomson scattering

Filtered Thomson Scattering: n_e , T_e , and EEDF inference

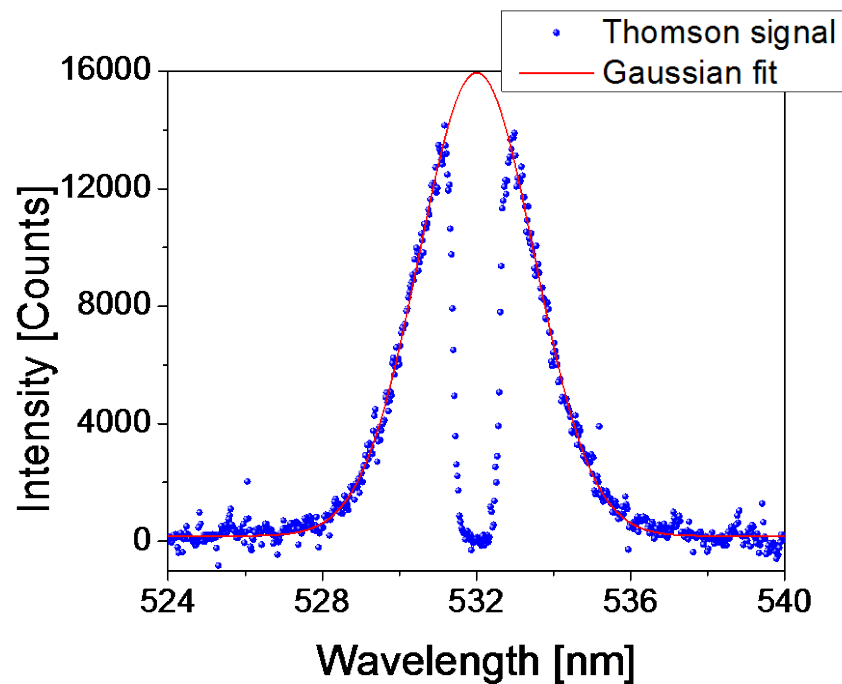
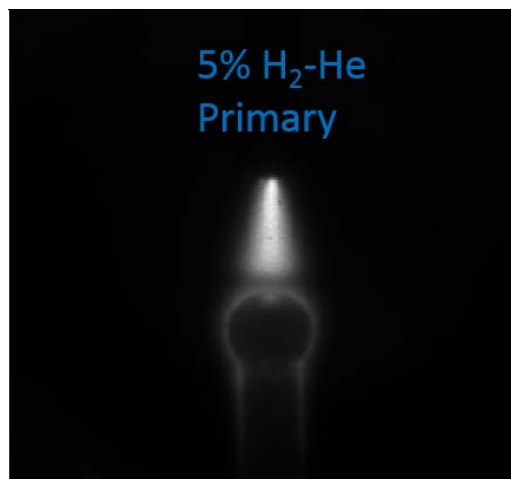


- Electron density: area under Thomson scattering spectrum
- Electron temperature: spectral linewidth
- Raman scattering rotational transitions in N_2 used for absolute calibration
- Gaussian Thomson scattering lineshape: Maxwellian EEDF

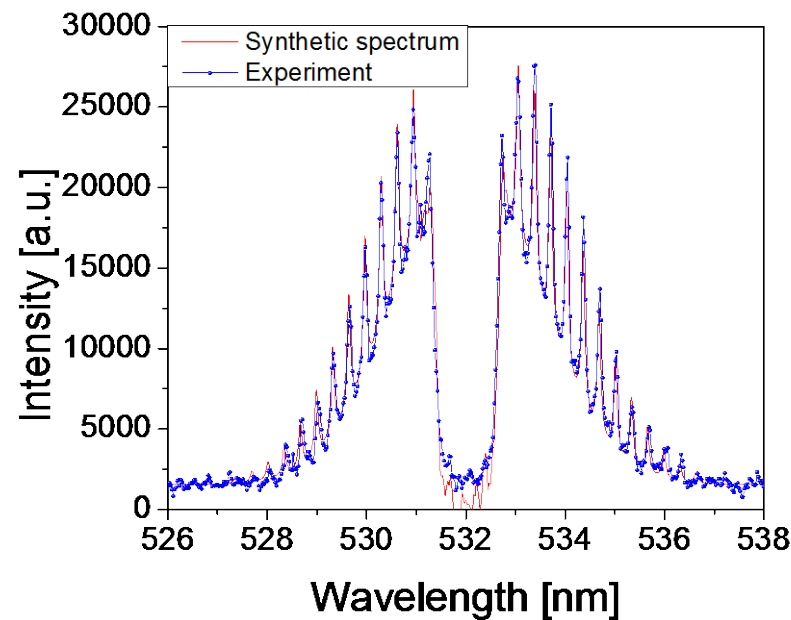
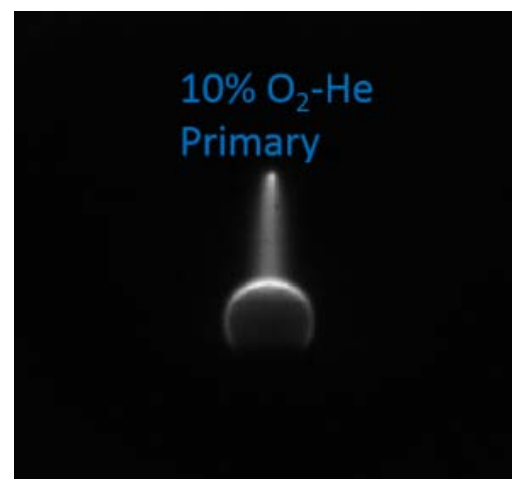


Thomson Scattering Spectra

Sphere-to-sphere ns pulse discharge in H₂-He and O₂-He



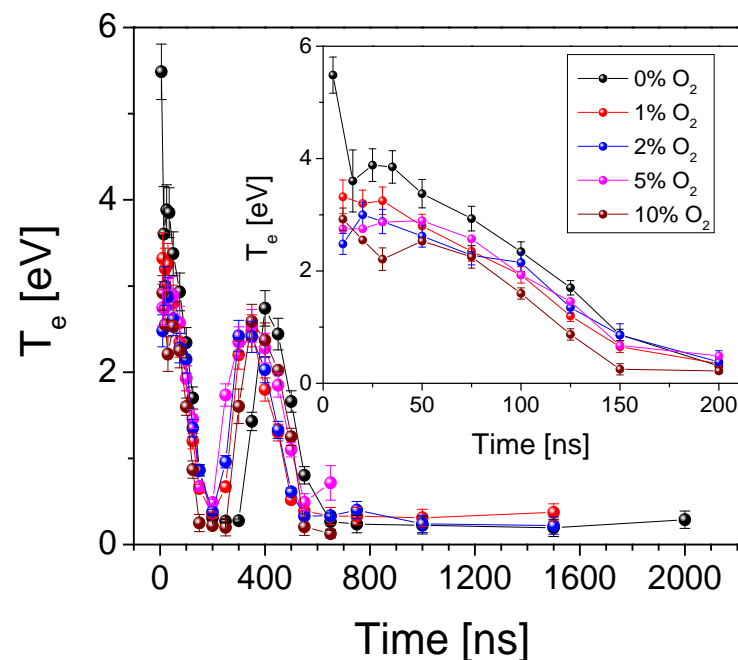
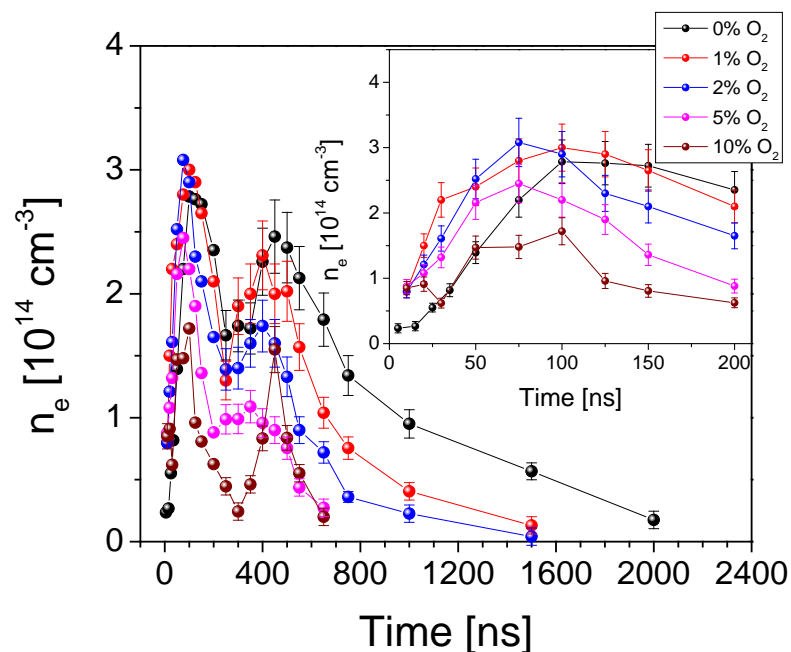
5% H₂-He, P=100 torr, t=100 ns
 $n_e = 1.5 \cdot 10^{14} \text{ cm}^{-3}$, $T_e = 2.0 \text{ eV}$



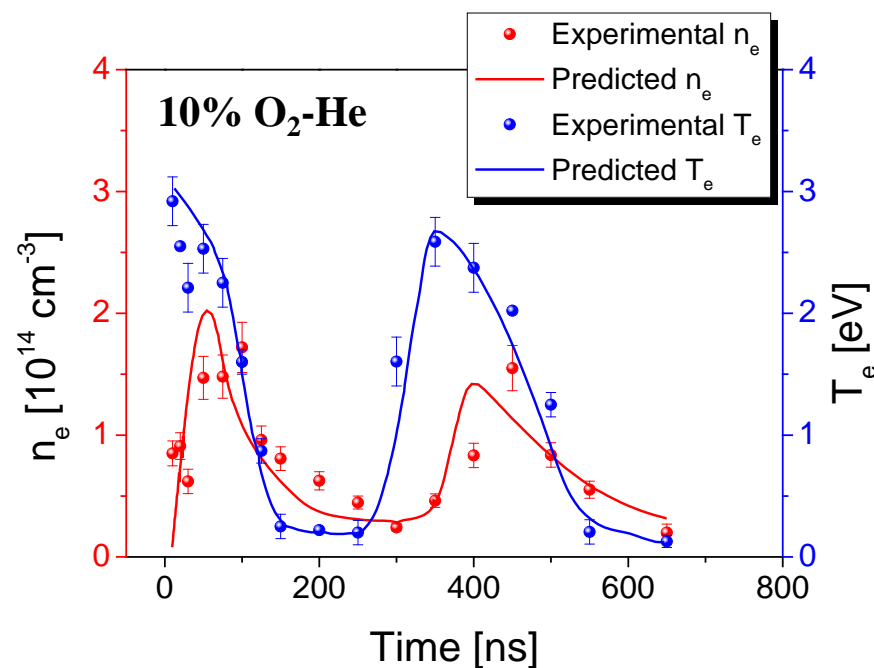
10% O₂-He, P=100 torr, t=100 ns
 $n_e = 1.7 \cdot 10^{13} \text{ cm}^{-3}$, $T_e = 1.6 \text{ eV}$, T=350 K

Electron Density and Electron Temperature

Sphere-to-sphere ns pulse discharge in O₂-He



- $n_e = 10^{13} - 3 \cdot 10^{14} \text{ cm}^{-3}$, $T_e = 0.3 - 5.5 \text{ eV}$ (0-10% O₂)
- “Double maxima” in n_e , T_e : two discharge pulses $\approx 400 \text{ ns}$ apart
- Electron temperature in the afterglow $T_e \approx 0.3 \text{ eV}$
Superelastic collisions prevent electron cooling
- Modeling predictions in good agreement with data

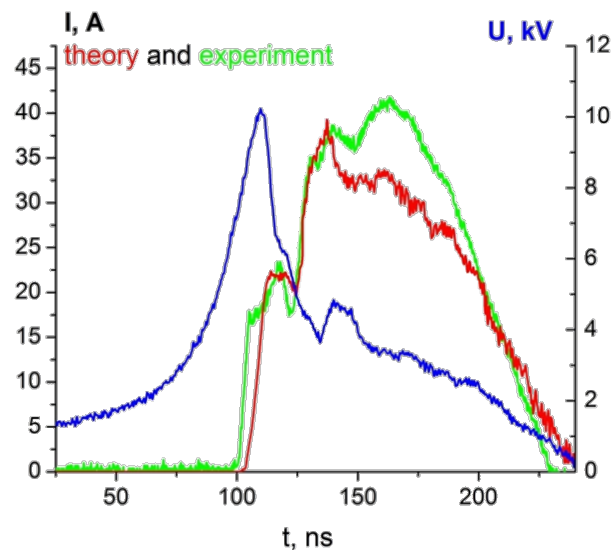
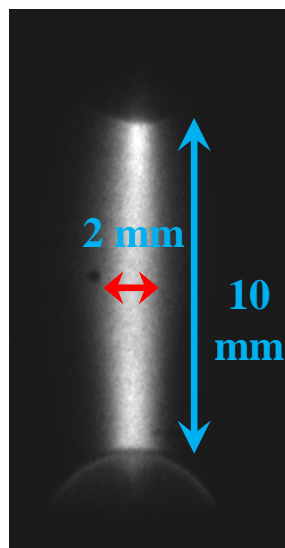


IV. Dynamics of temperature rise in transient plasmas: “rapid” heating and “slow” heating

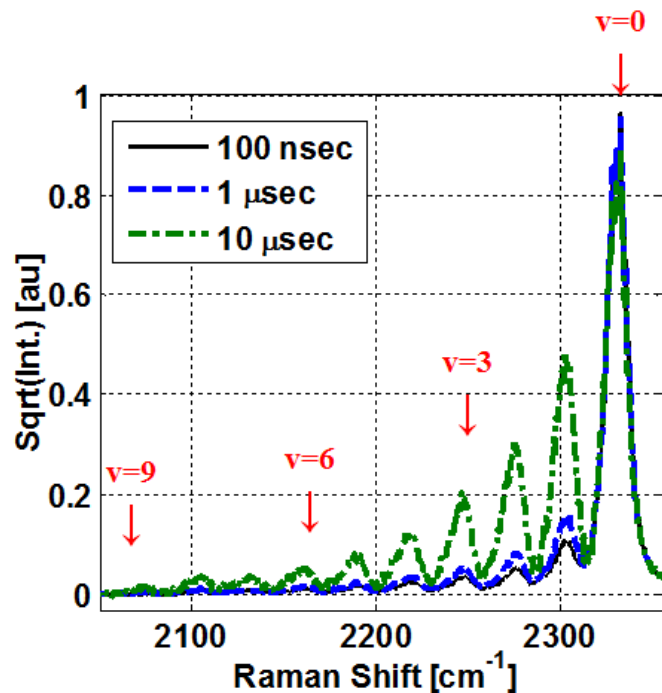
Diagnostics: vibrational and pure rotational ps CARS

Sphere-to-sphere ns pulse discharge in air, P=100 Torr

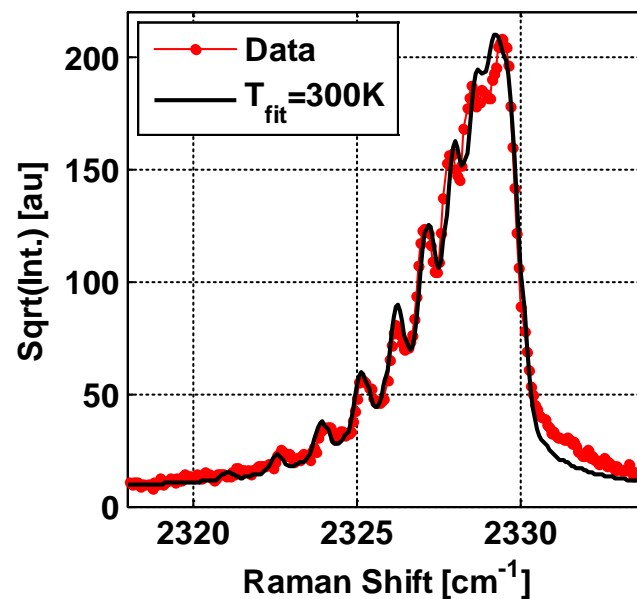
Discharge pulse waveforms and CARS spectra



Pulse energy coupled to plasma ~ 0.5 eV



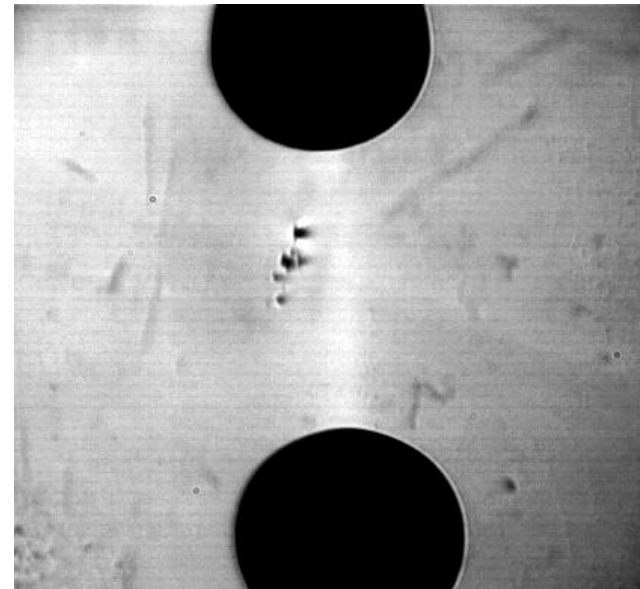
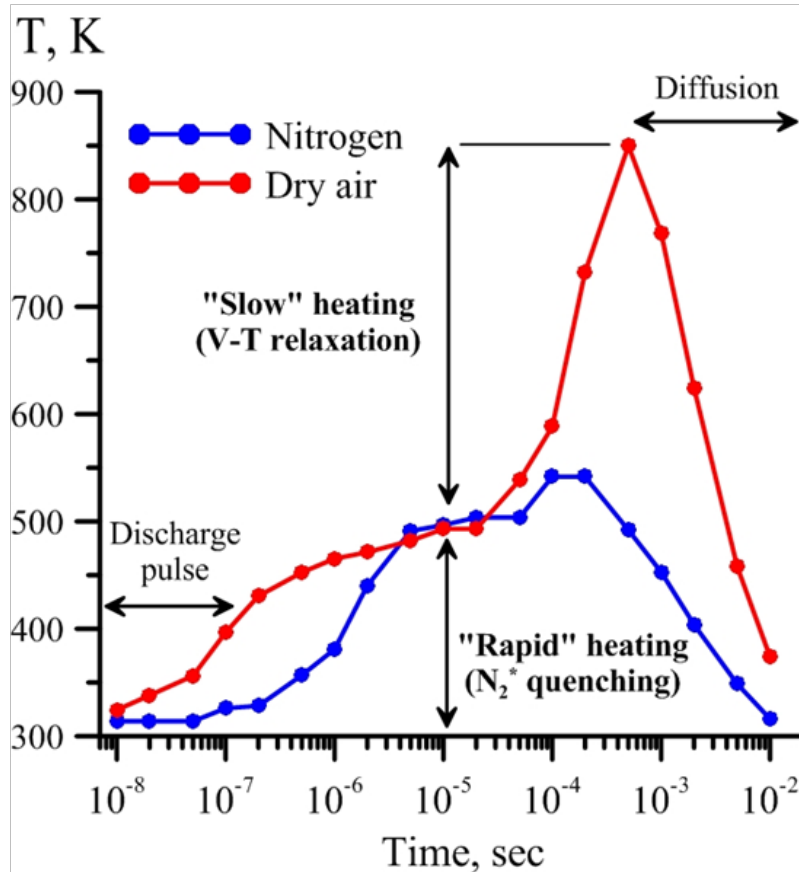
$N_2(v=0-9)$ bands during and after discharge pulse



$N_2(v=0)$ band (without plasma)

Temperature rise in ns pulse discharge and afterglow: air vs. nitrogen

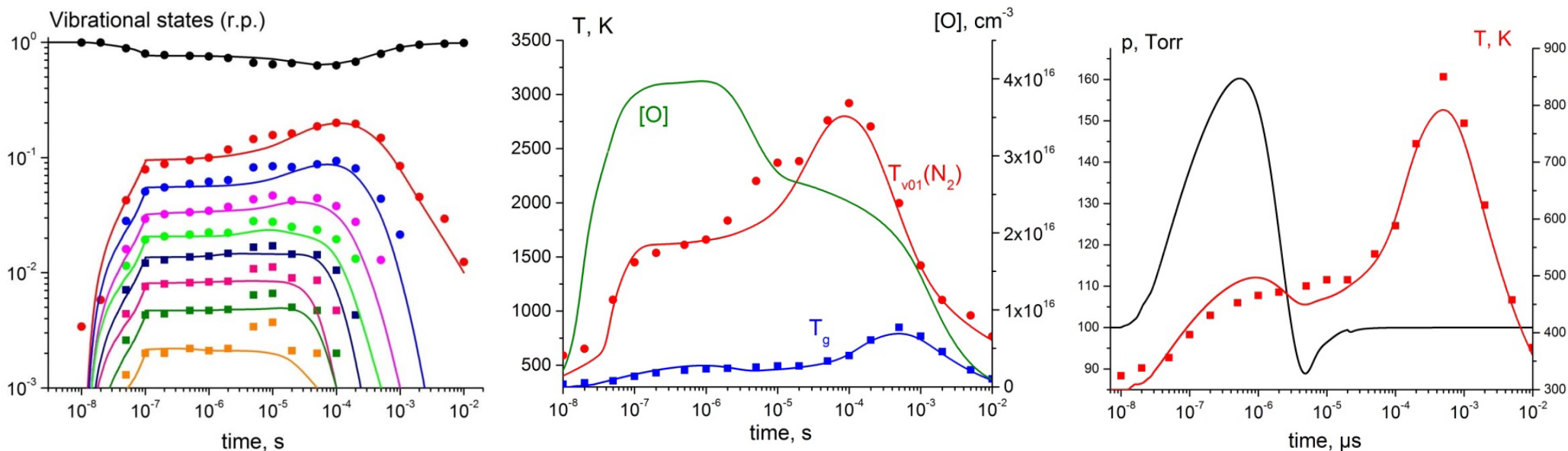
Air, P=100 Torr



$t = 1-10 \mu\text{s}$ (frames are $1 \mu\text{s}$ apart)

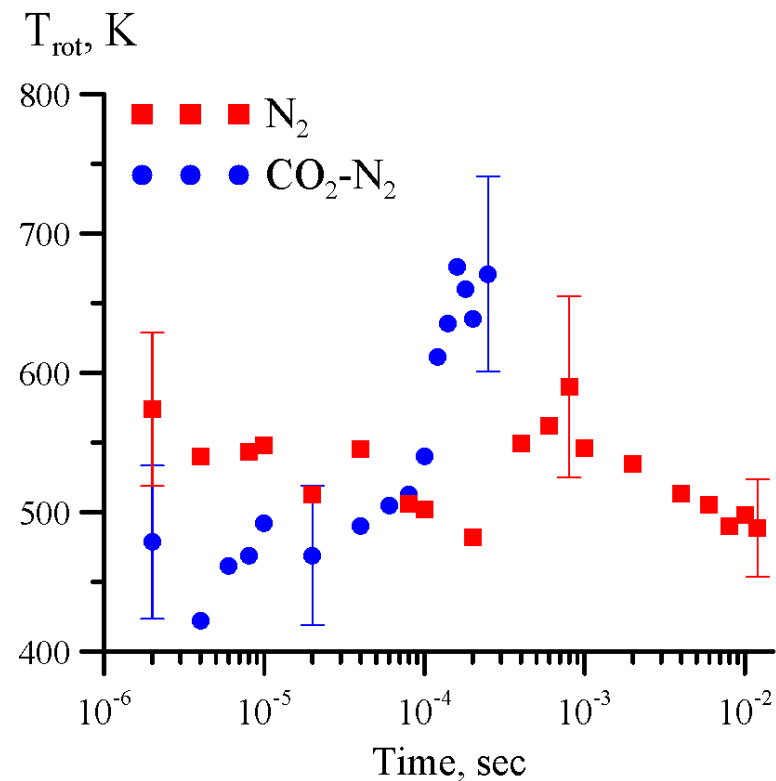
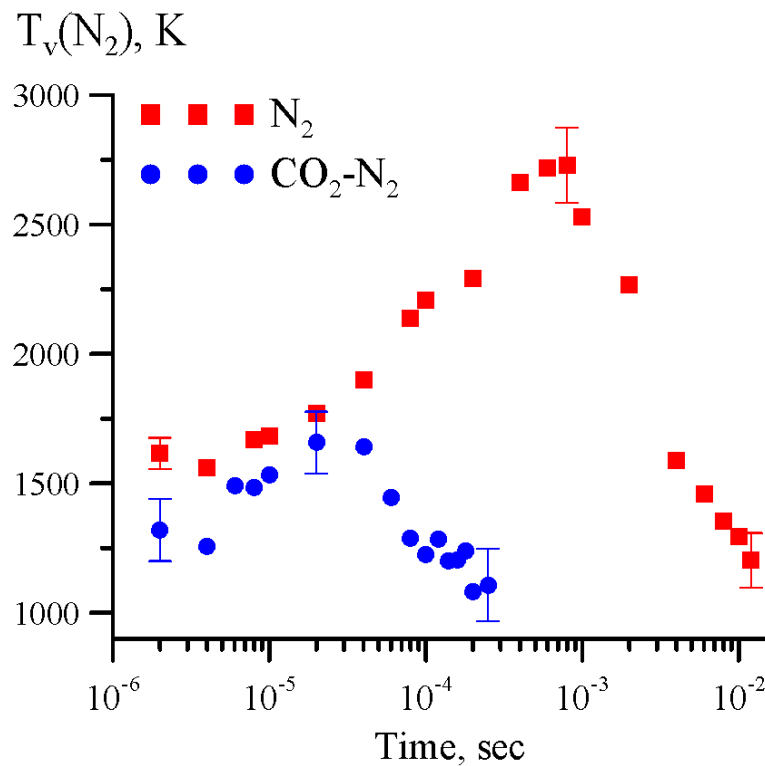
- Compression waves formed by "rapid" heating, on sub-acoustic time scale, $\tau_{\text{acoustic}} \sim r / a \sim 2 \mu\text{s}$
- Strong effect on high-speed flows

Comparison with modeling predictions in air: vibrational kinetics and temperature rise



- Strong vibrational excitation in the discharge, $N_2(v=0-8)$
- $T_v(N_2)$ rise in early afterglow: V-V exchange, $N_2(v) + N_2(v=0) \rightarrow N_2(v-1) + N_2(v=1)$
- $T_v(N_2)$ decay in late afterglow: V-T relaxation, $N_2(v) + O \rightarrow N_2(v-1) + O$, radial diffusion
- “Rapid” heating: quenching of N_2 electronic states, $N_2(C,B,A,a) + O_2 \rightarrow N_2(X) + O + O$
- “Slow” heating: V-T relaxation, $N_2(X,v) + O \rightarrow N_2(X,v-1) + O$
- “Rapid” heating: pressure overshoot on centerline, compression waves detected in experiments
- NO formation: dominated by reactions of N_2 electronic states, $N_2^* + O \rightarrow NO + N$, also in good agreement with [NO], [N] measurements

Adding CO₂ (rapid V-T relaxer) to air: accelerating energy thermalization rate



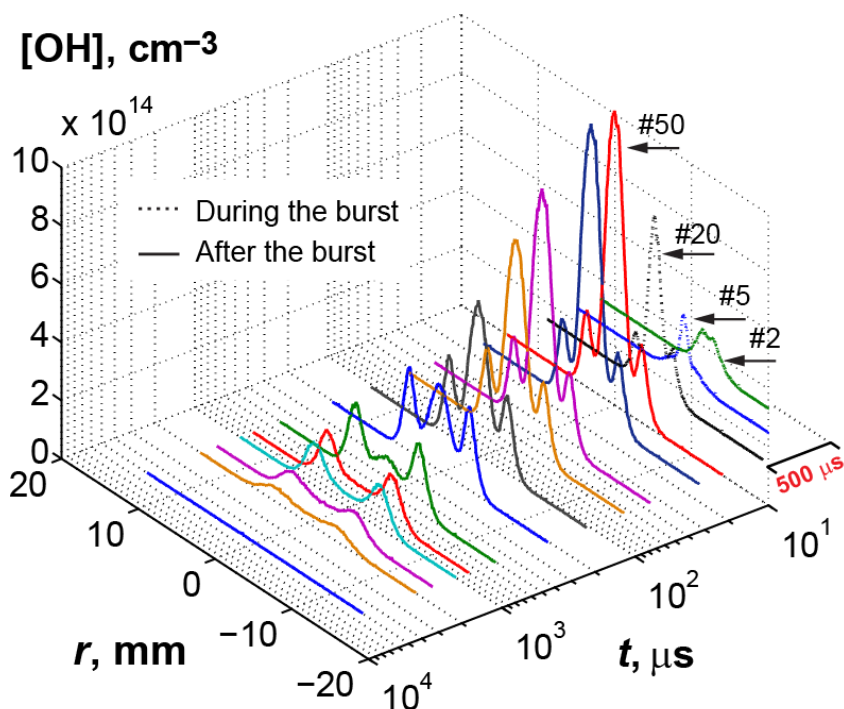
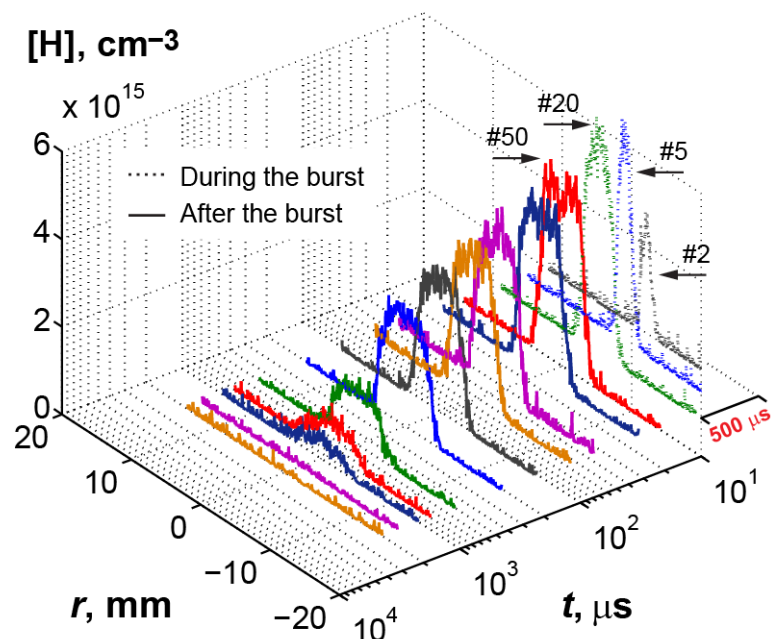
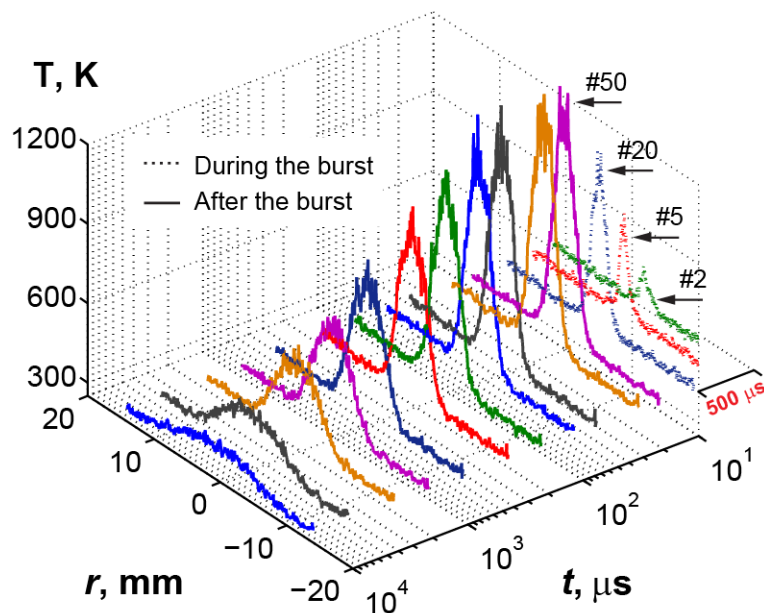
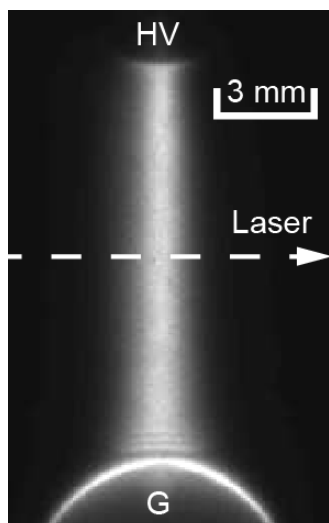
Rapid N₂ relaxation and temperature rise. Mechanism of accelerated heating:

- V-V energy exchange between N₂ and CO₂(v₃) mode: $N_2(v=1) + CO_2(000) \leftrightarrow N_2(v=0) + CO_2(001)$
- CO₂ energy re-distribution among vibrational modes: $CO_2(001) + M \leftrightarrow CO_2(100,020,010) + M$
- V-T relaxation of bending mode: $CO_2(010) + M \rightarrow CO_2(100) + M$
- Strong effect on nonequilibrium compressible flows

**V. Fuel-air chemistry in transient plasmas:
kinetics of plasma assisted combustion**

Diagnostics: Rayleigh scattering, LIF, CARS

Plasma chemical reactions and transport: ns pulse discharge in $\text{H}_2 - \text{O}_2 - \text{Ar}$, $P=40$ torr



Hot central region:

OH production dominated by chain branching



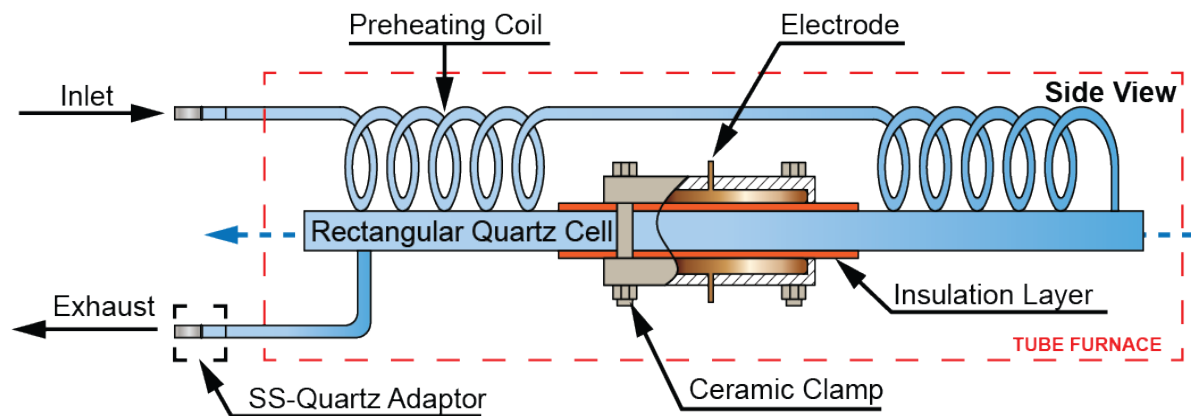
Colder peripheral region:

OH accumulation in HO_2 reactions

radial diffusion of H atoms;



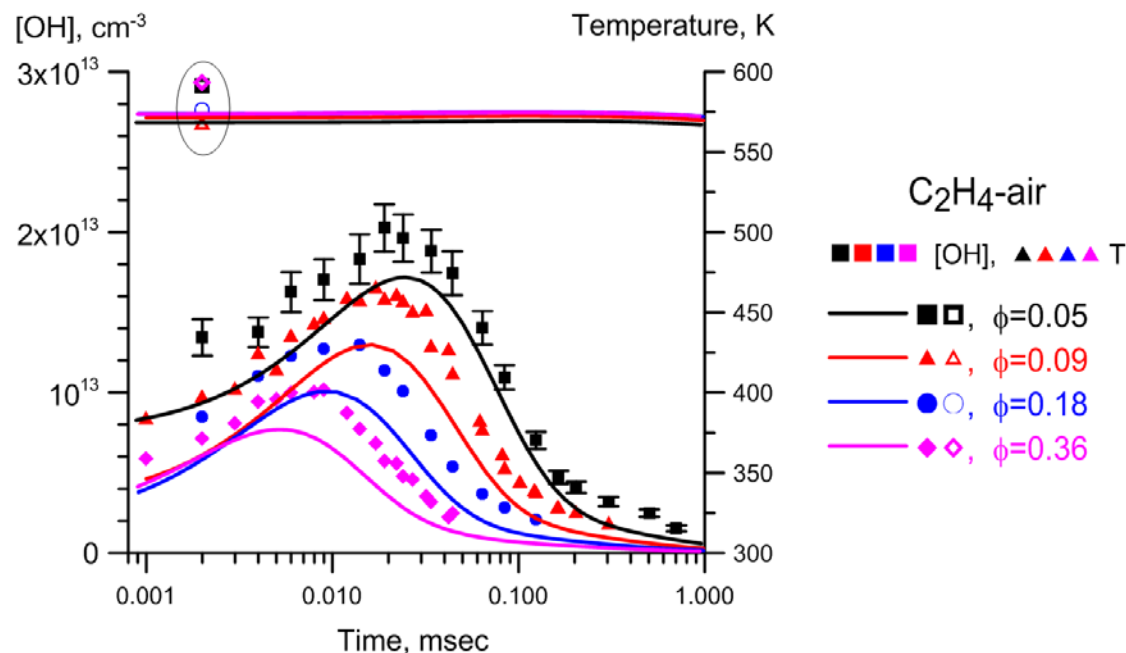
[OH] kinetics in preheated fuel-air mixtures after ns pulse discharge burst



End View

	Pulse #10	Pulse #100
$\text{H}_2 - \text{air}, \phi=0.3$ $T_0=500 \text{ K}, P=100 \text{ torr}$		
$\text{C}_2\text{H}_4 - \text{air}, \phi=0.3$ $T_0=500 \text{ K}, P=100 \text{ torr}$		

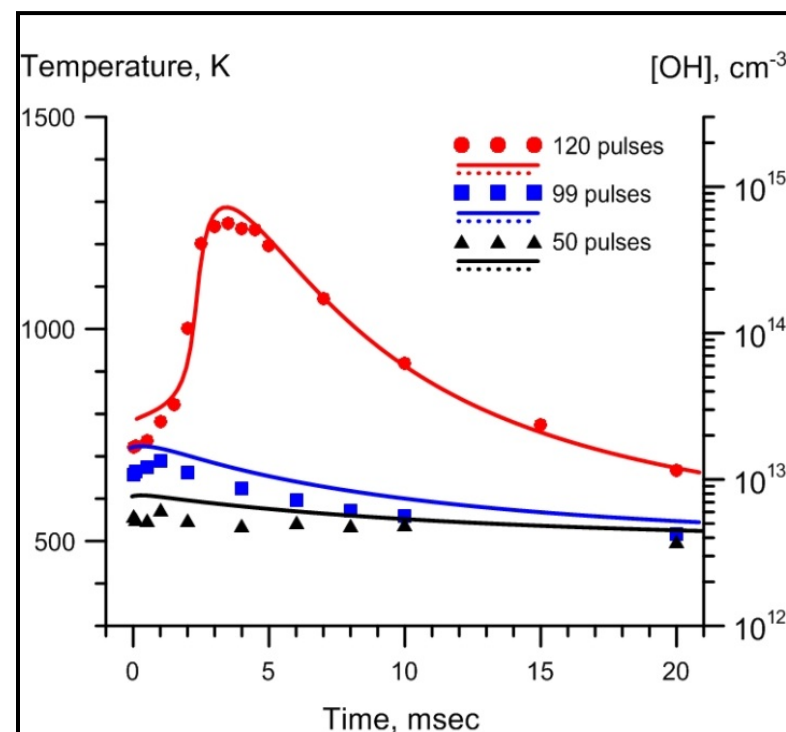
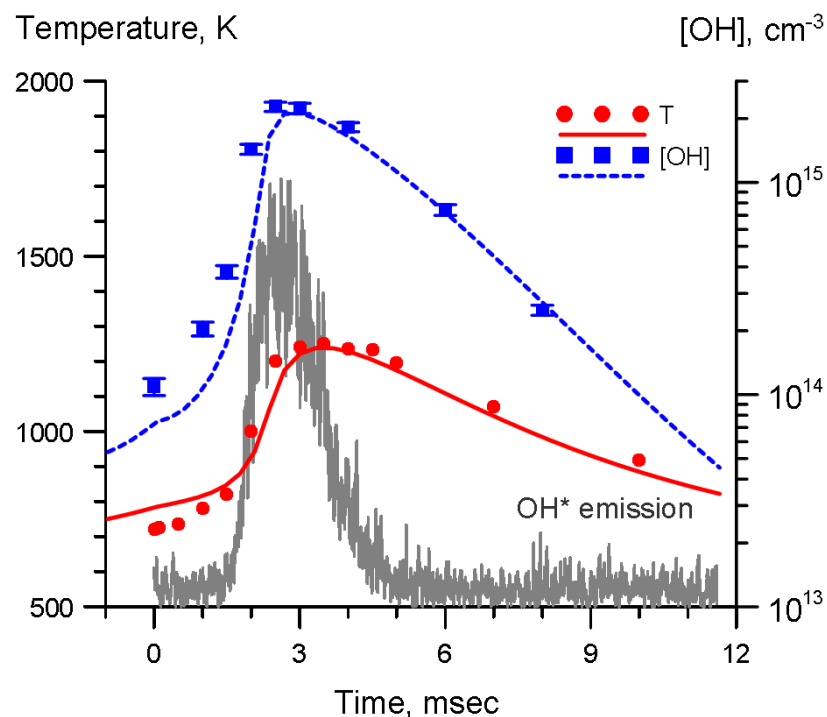
“Near 0-D” diffuse plasma, a burst of 50-100 pulses used to couple sufficient energy



Comparison with modeling predictions: plasma assisted combustion kinetic mechanism validation

Plasma chemical reactions

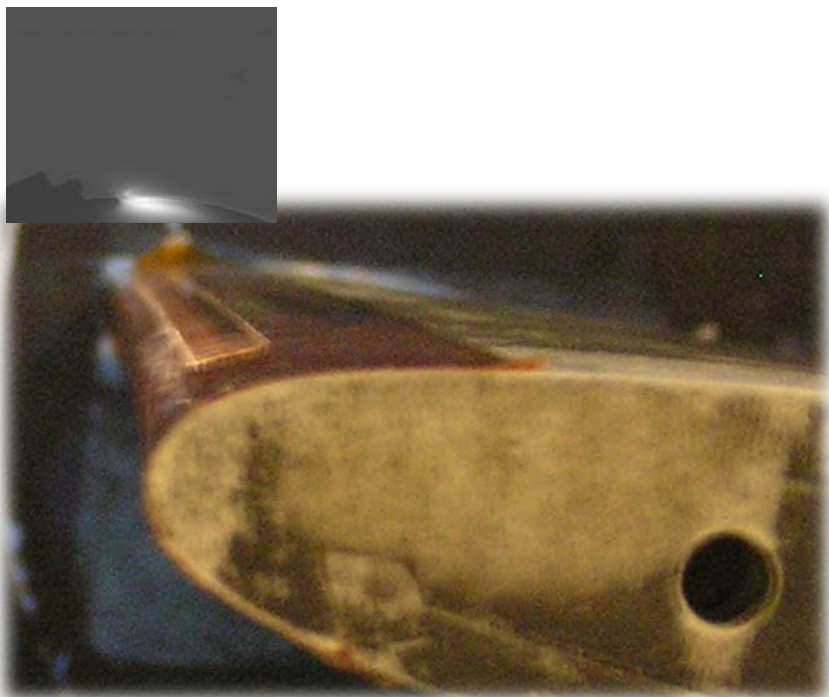
reduce ignition temperature in H₂-air



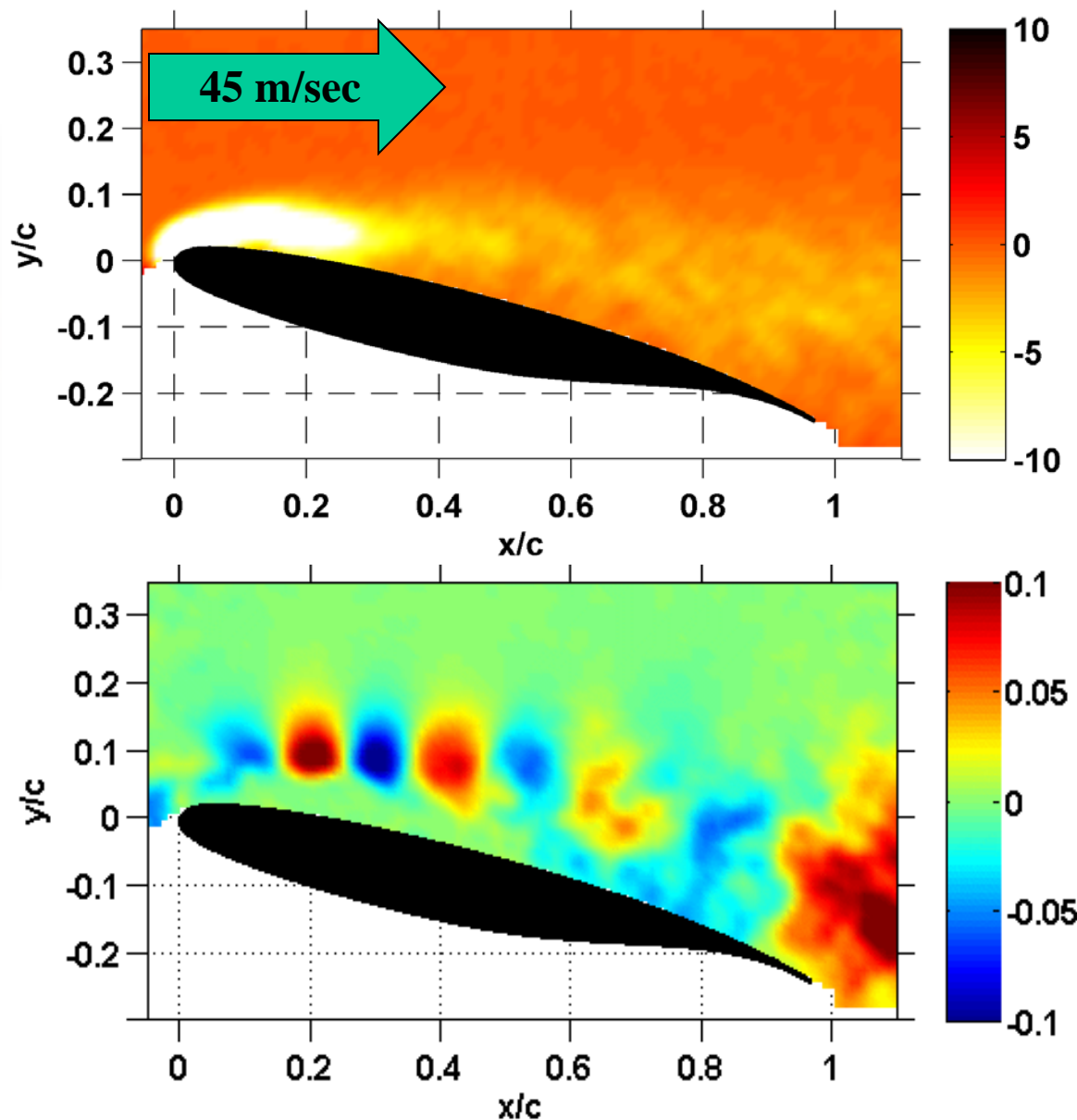
- $\phi=0.4$, 120-pulse burst, $T_0=500$ K, $P=80-90$ torr
- Model predictions: good agreement with time-resolved temperature measurements
- Ignition temperature with plasma, $T_i \approx 700$ K, lower than autoignition temperature, $T_a \approx 900$ K

VI. Air plasma kinetics and plasma flow control

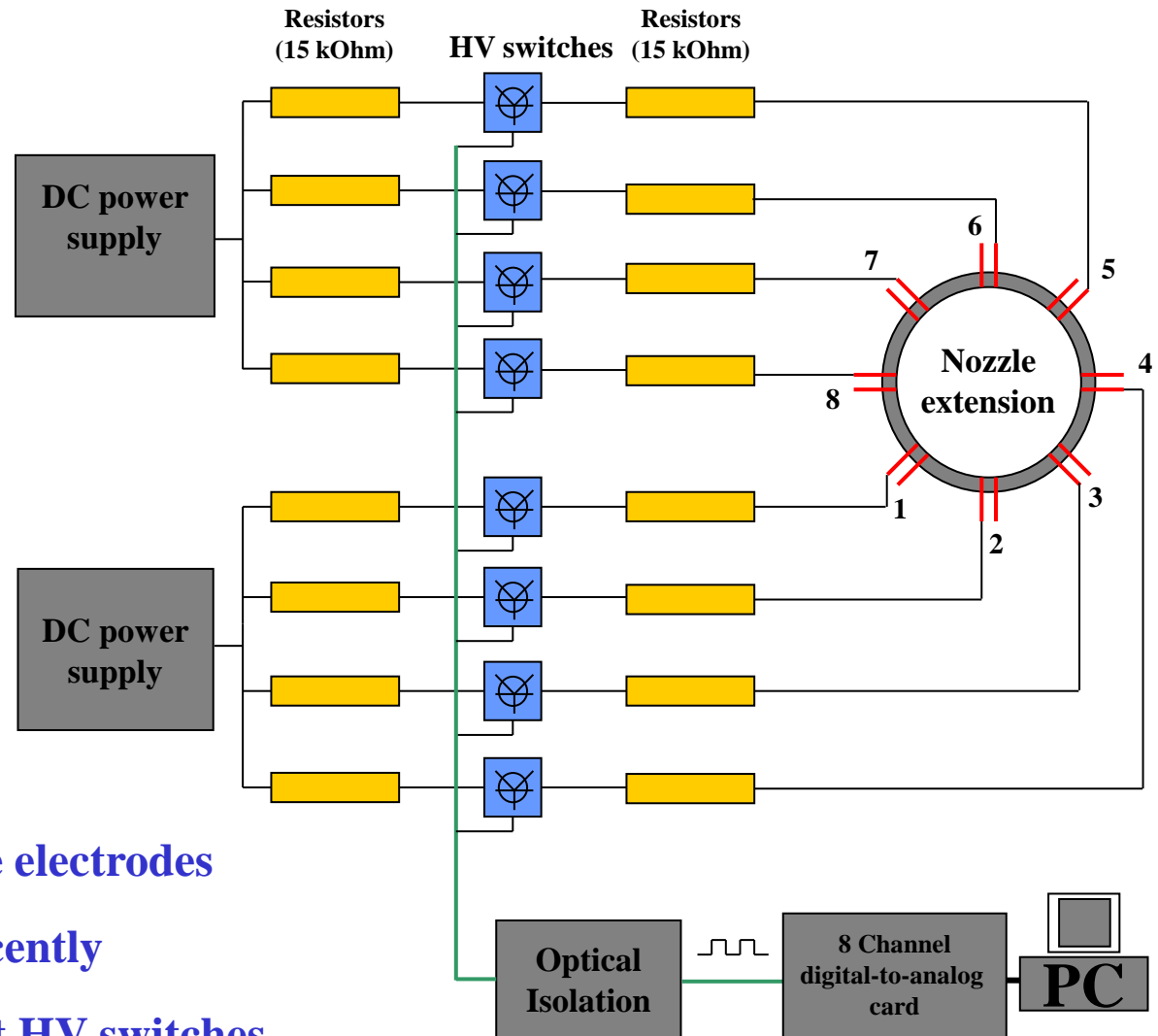
Ns DBD plasma actuators: subsonic flow boundary layer reattachment



- Every nanosecond discharge pulse produces a robust spanwise vortex
- Enhanced mixing with free stream \rightarrow boundary layer reattachment
- Same effect detected up to $u=96$ m/sec ($M=0.28$, $Re_x \sim 1.5 \cdot 10^6$)
- Consistently outperform AC DBD actuators

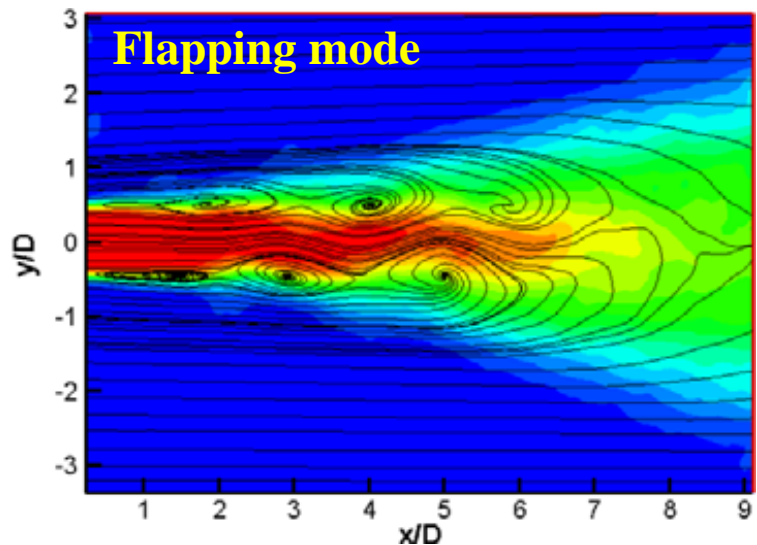
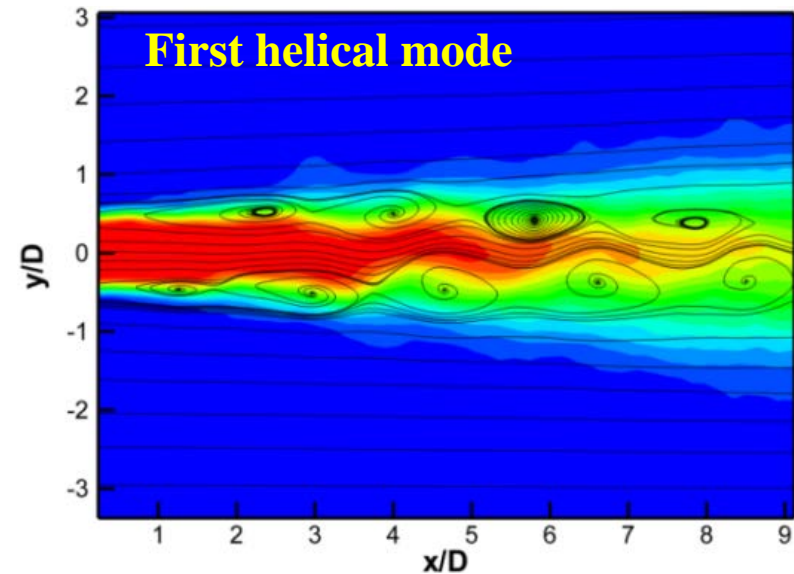
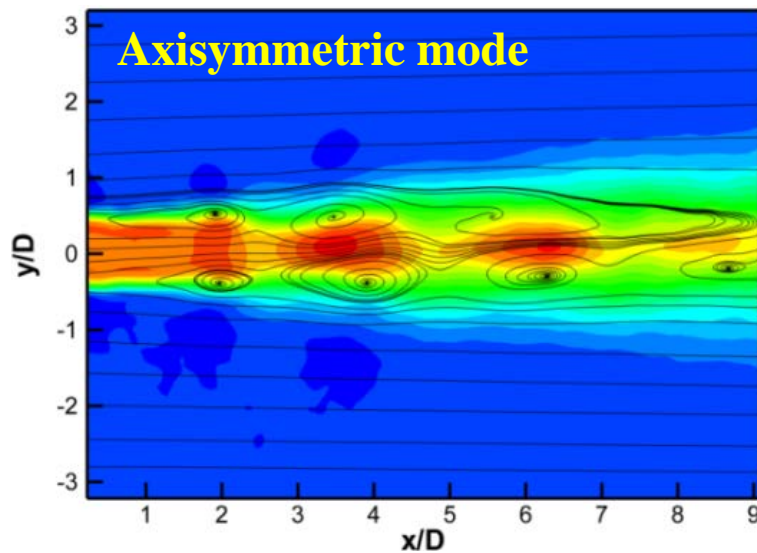


Localized Arc Plasma Flow Actuators (LAPFA): Exciting instabilities in transonic and supersonic flows ($M=0.9-2.0$)



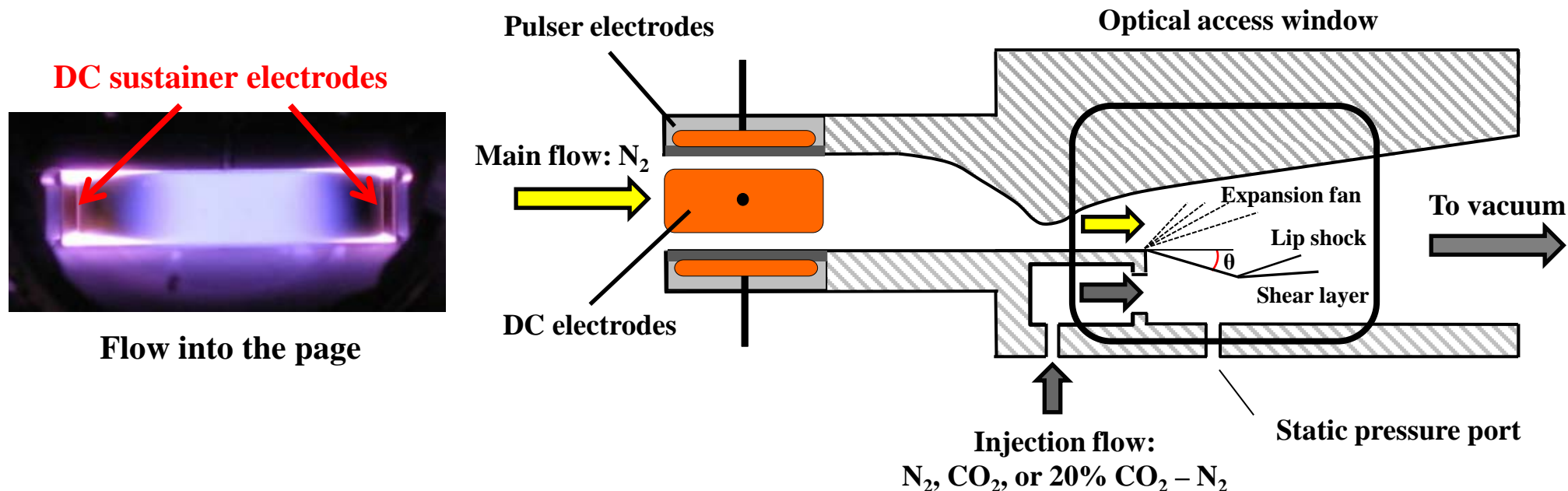
- BN nozzle extensions, tungsten wire electrodes
- Circular nozzle, 1 inch diameter recently
- Multiple channels controlled by fast HV switches
- Independent control of frequency, phase, and duty cycle → excitation of different instability modes

LAPFA: Formation of coherent structures in a $M=0.9$ circular jet



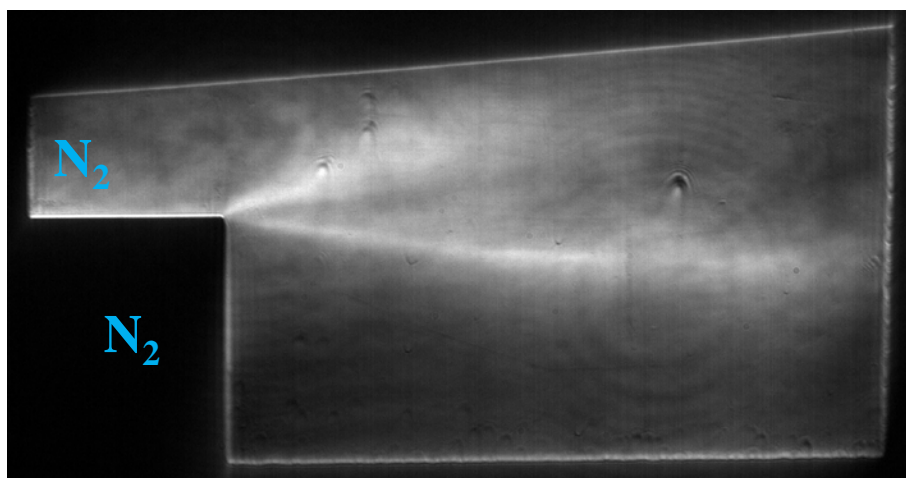
- High amplitude perturbations (localized heating in arc filaments)
- Every discharge pulse results in vortex formation
- Flow responds to forcing near jet column instability frequency

Control of supersonic mixing / shear layer by accelerated relaxation of vibrational energy

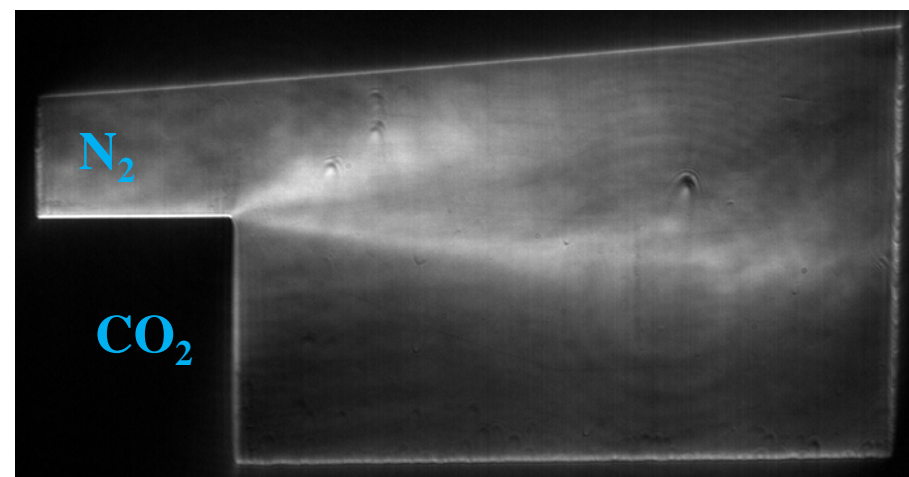


- Plenum: overlapped ns pulse / DC sustainer discharge for vibrational loading of N_2
- $P_0 = 300$ torr, $T_v = 2000$, $T = 500$ K, 2-D nozzle, top wall contoured, bottom wall plane
- Condition at nozzle exit : $M = 2.5$, $P_{exit} = 15$ torr
- Subsonic flow below expansion corner: injection of N_2 or CO_2
- Optical access for schlieren, CARS, and NO PLIF in subsonic and supersonic flows

Effect of vibrational relaxation of shear layer: N_2 / N_2 (left) vs. N_2 / CO_2 (right)



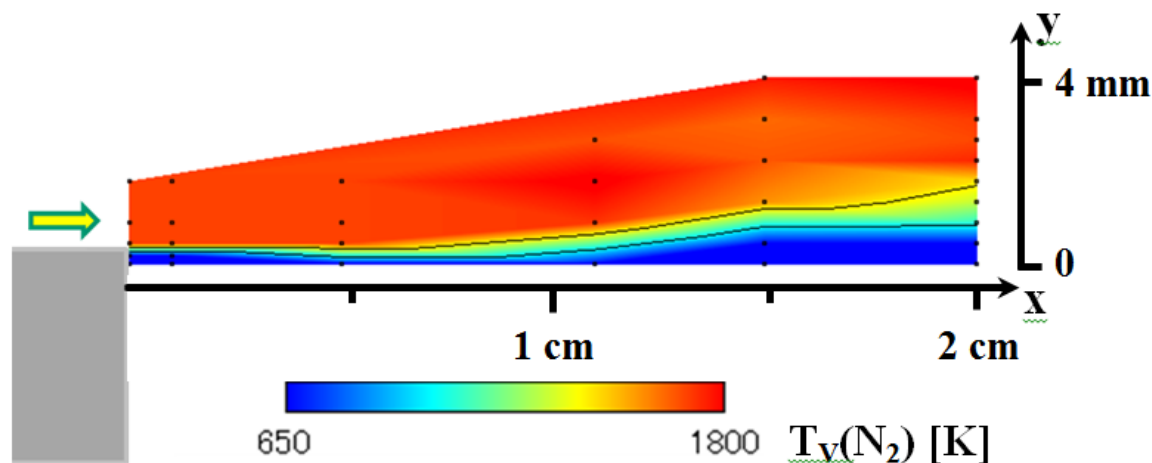
N_2 “bleeding” through backstep



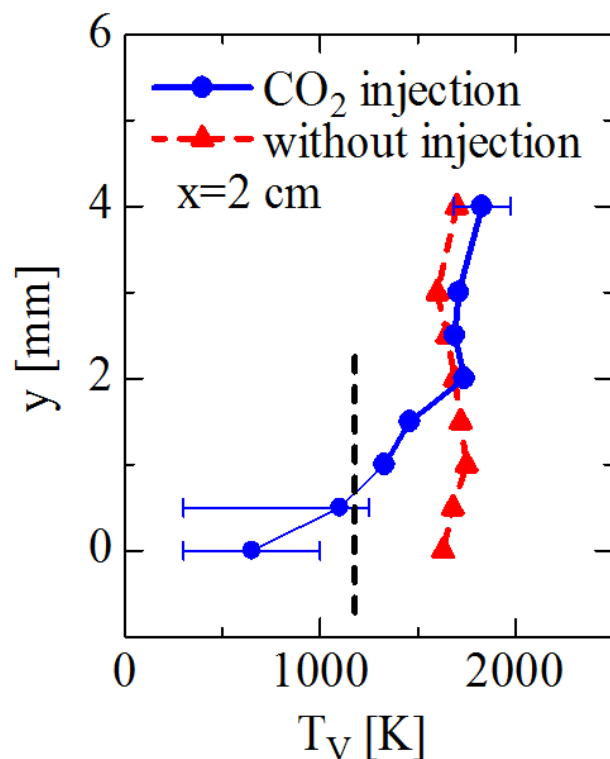
CO_2 “bleeding” through backstep

- Time delay between frames 5 ms, $t = 0-80$ ms
- N_s pulse / DC discharge (2.3 kW) is turned on at $t = 10-45$ ms, to excite main N_2 flow
- No perturbation of shear layer detected in N_2 / N_2 flow
- In N_2 / CO_2 flow, shear layer expansion angle decreases, approaching $\theta=0^\circ$
- No change observed if main N_2 flow is not excited

N₂ Vibrational Temperature Distribution in Shear Layer



- Top flow: vibrationally excited N₂, T_v=1900 K, estimated T_{rot}=240 K
- Bottom flow: CO₂ bleeding through backstep, static pressure 7 torr



- CO₂ bleeding reduces T_v(N₂), increases T_{trans/rot} and static pressure
- Consistent with time-resolved measurements in ns pulse discharge in quiescent N₂-CO₂
- Static pressure increase pushes up shear / mixing layer

Summary: air plasma kinetics

- Growing body of time-resolved, spatially-resolved data characterizing transient, high-pressure air and fuel-air plasmas
- Measurements of electric field, electron density, and electron temperature necessary for insight into discharge energy coupling and partition
- Measurements of temperature, $N_2(v)$ populations, and excited electronic states of N_2^* necessary for insight into temperature dynamics
- Measurements of N_2^* and key radicals (O, H, OH, and NO) critical for quantifying their effect on fuel-air plasma chemistry
- Comparing measurement results with kinetic modeling predictions provides confidence in the models, assesses their predictive capability

Summary: nonequilibrium plasma flow control

- **Surface and volumetric ns pulse discharges: energy thermalization on sub-acoustic time scale, high-amplitude compression wave generation**
- **Mechanism of energy thermalization (“rapid heating” and “slow heating”) is well understood**
- **NS-DBD surface plasma actuators: large-scale coherent flow structures; significant flow control authority in subsonic flows (up to $M = 0.3$) at low actuator powers; scalable to large dimensions (~1 m)**
- **LAFPA actuators: large-scale coherent structures; excitation of flow instability modes; significant control authority in transonic and supersonic flows ($M = 0.9$ - 2.0) at low actuator powers; scalable to large phased arrays**
- **Flow control by vibrational relaxation: injection of “rapid relaxer” species into nonequilibrium flow at desired location; temperature and pressure rise due to accelerated relaxation; strong effect in supersonic shear layer**

Acknowledgments

AFOSR MURI “Fundamental Mechanisms, Predictive Modeling, and Novel Aerospace Applications of Plasma Assisted Combustion”

AFOSR BRI “Nonequilibrium Molecular Energy Coupling and Conversion Mechanisms for Efficient Control of High-Speed Flow Fields”

DOE PSAAP-2 Center “Exascale Simulation of Plasma-Coupled Combustion” (under U. Illinois at Urbana-Champaign prime)

US DOE Plasma Science Center “Predictive Control of Plasma Kinetics: Multi-Phase and Bounded Systems”

NSF “Fundamental Studies of Accelerated Low Temperature Combustion Kinetics by Nonequilibrium Plasmas”