

Electric Field Measurements in Atmospheric Pressure Electric Discharges

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Acknowledgments

**US DOE Plasma Science Center “Predictive Control of Plasma
Kinetics: Multi-Phase and Bounded Systems”**

**NSF “Fundamental Studies of Accelerated Low Temperature
Combustion Kinetics by Nonequilibrium Plasmas”**

Motivation

- **High-pressure, low-temperature air plasmas:** applications for plasma-assisted combustion, plasma flow control, plasma surface processing, biology and medicine
- **Ns pulse discharges in air and fuel-air mixtures:** stable at high pressures, efficient generation of excited species and reactive radicals
- **Electric field in plasmas:** controlled by ionization, electron and ion transport; electron emission from electrodes; surface charge accumulation on dielectrics
- **Electric field controls energy partition in the plasma** (vibrational and electronic excitation, dissociation, ionization), rate of gas temperature rise
- **Need for non-intrusive, spatially resolved, time-resolved electric field measurements** in transient reacting plasmas, near surfaces
- **Insight into kinetics of ionization and charge transport, validation of kinetic models**

Energy Partition in Air Plasma vs. Electric Field

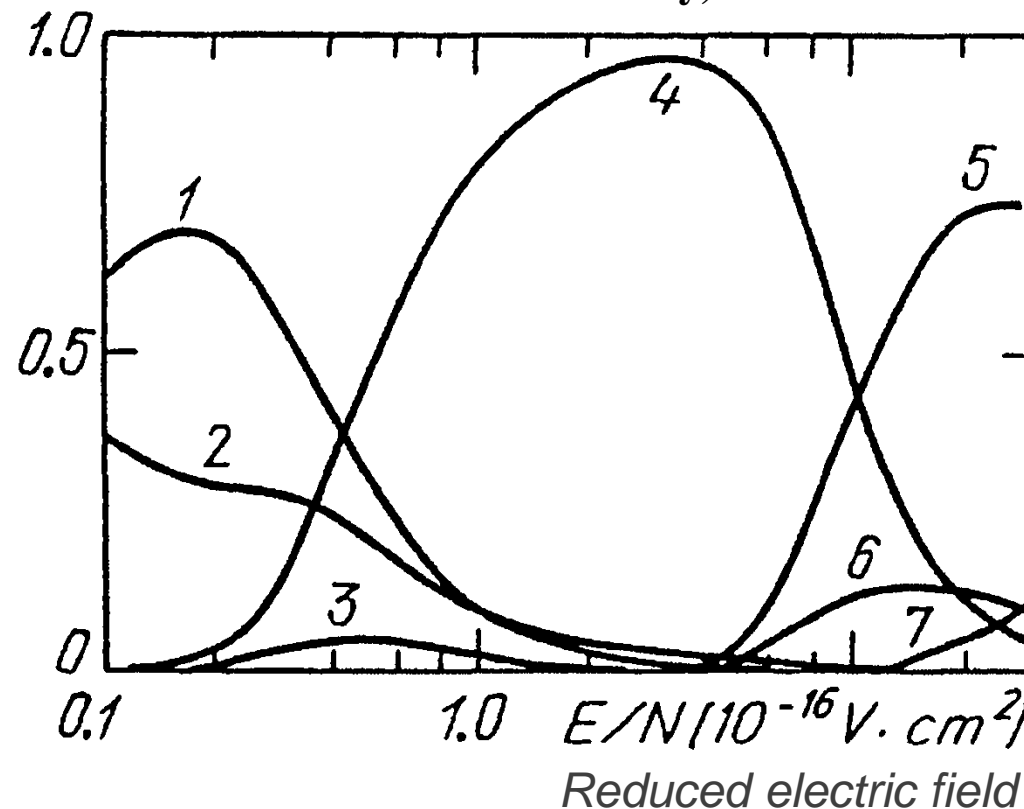
Yu. Raizer, Gas Discharge Physics, Springer, 1991

Quasi-steady-state discharges, low E/N:

(4) **N₂ vibrational excitation**:

Low reactivity, slow thermalization

Energy fraction



Pulsed discharges, high E/N

(5,6) **electronic excitation**,
dissociation, (7) **ionization**

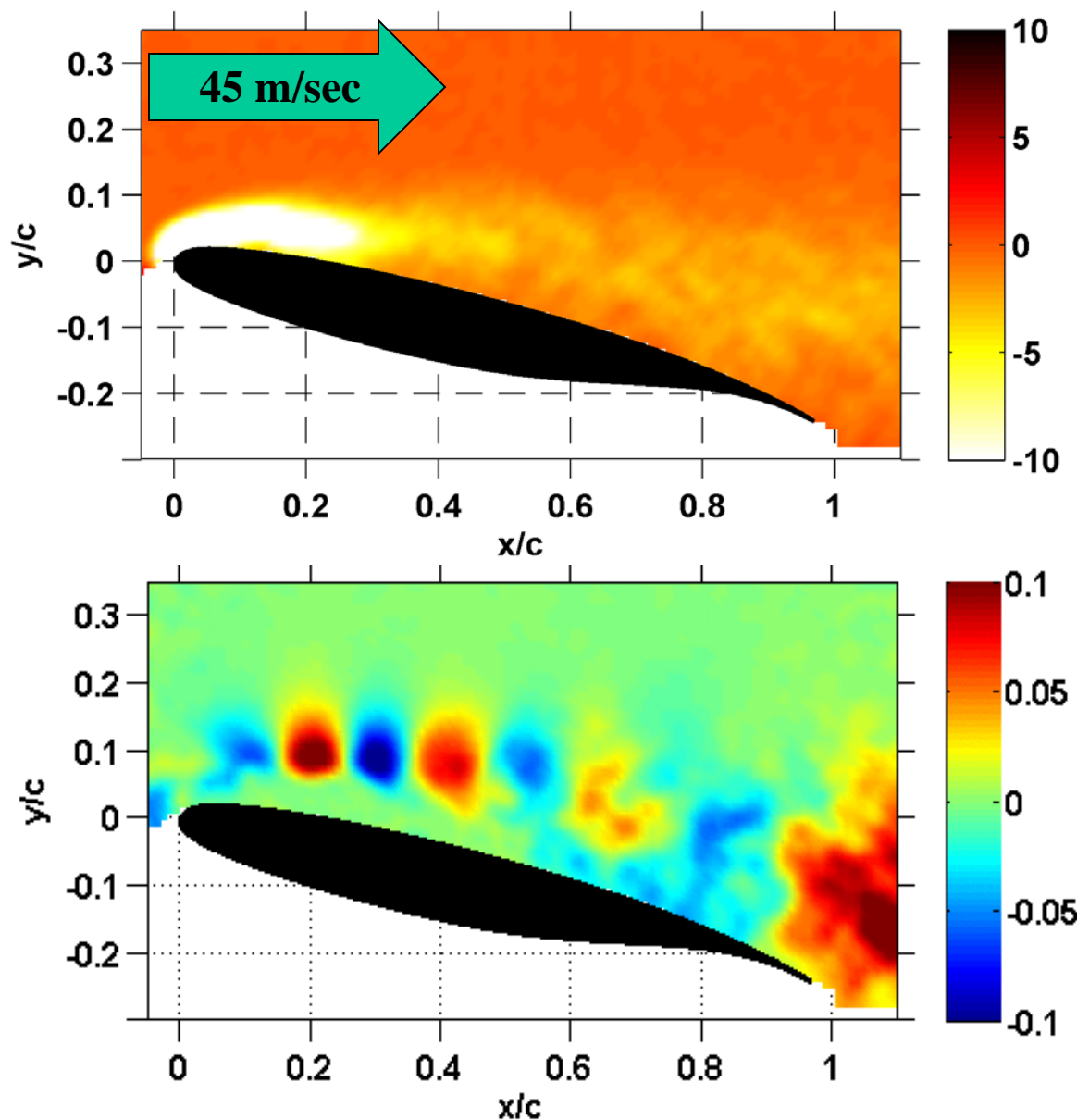
High reactivity, rapid thermalization

- Reduced electric field, E/N , controls input energy partition in the discharge
- Rates of electron impact processes: depend on E/N exponentially

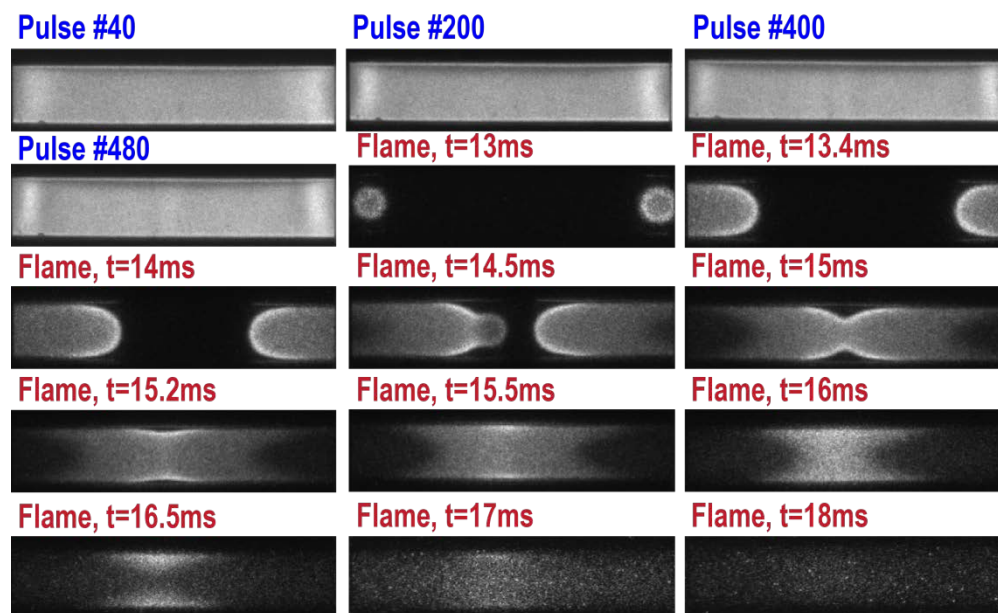
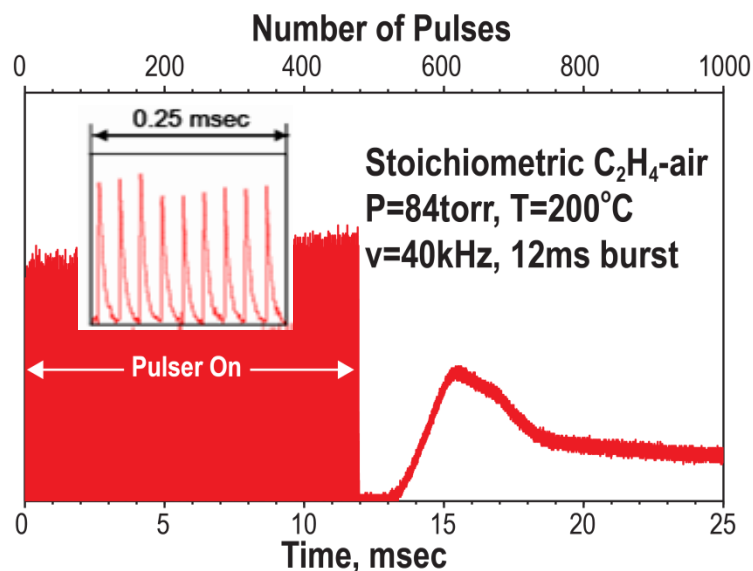
Ns Pulse Surface Plasma Actuators form spanwise vortices in flow over airfoil



- Every ns discharge pulse produces a spanwise vortex
- Enhanced mixing with free stream
→ **boundary layer reattachment**
- Effect detected up to $u=96$ m/sec
($M=0.28$, $Re_x \sim 1.5 \cdot 10^6$)



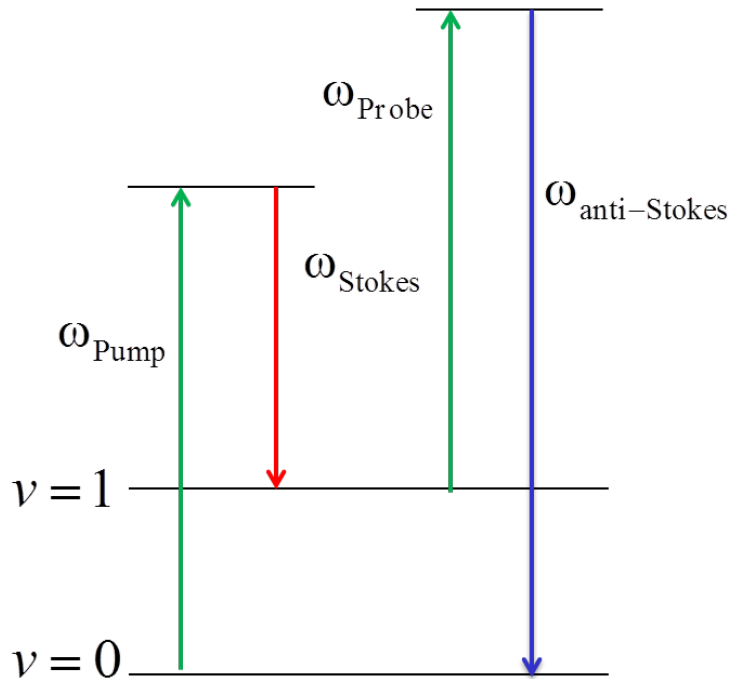
Pulse Burst Nanosecond Discharge: C_2H_4 -ignition below autoignition temperature



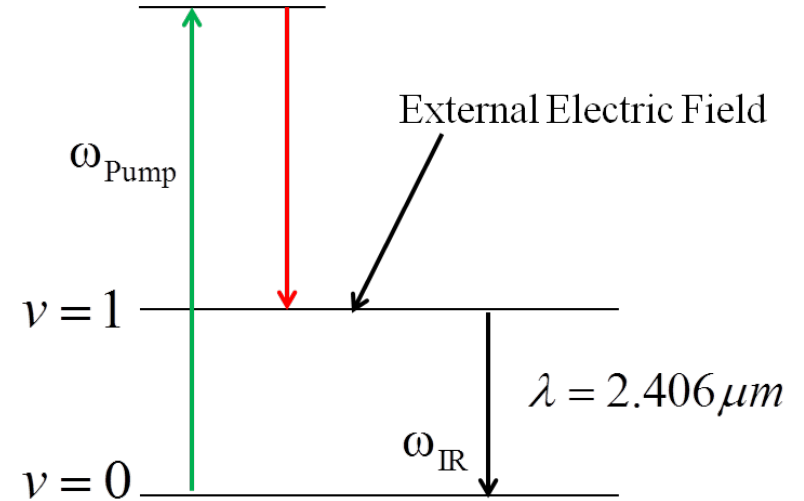
OH emission from plasma and flame

- Ignition induced by radicals generated in the plasma (primarily O and H atoms)
- Ignition occurs at temperature ≈ 200 K below autoignition
- Ignition begins near edges of the plasma (higher energy loading)
- Flame propagates to the center of the plasma

4-wave mixing: “CARS-like” process (E-field acting as a “zero-frequency” probe beam)



$$I_{\text{CARS}} \sim \chi_{\text{CARS}}^2 I_{\text{P}}^2 I_{\text{Stokes}}$$

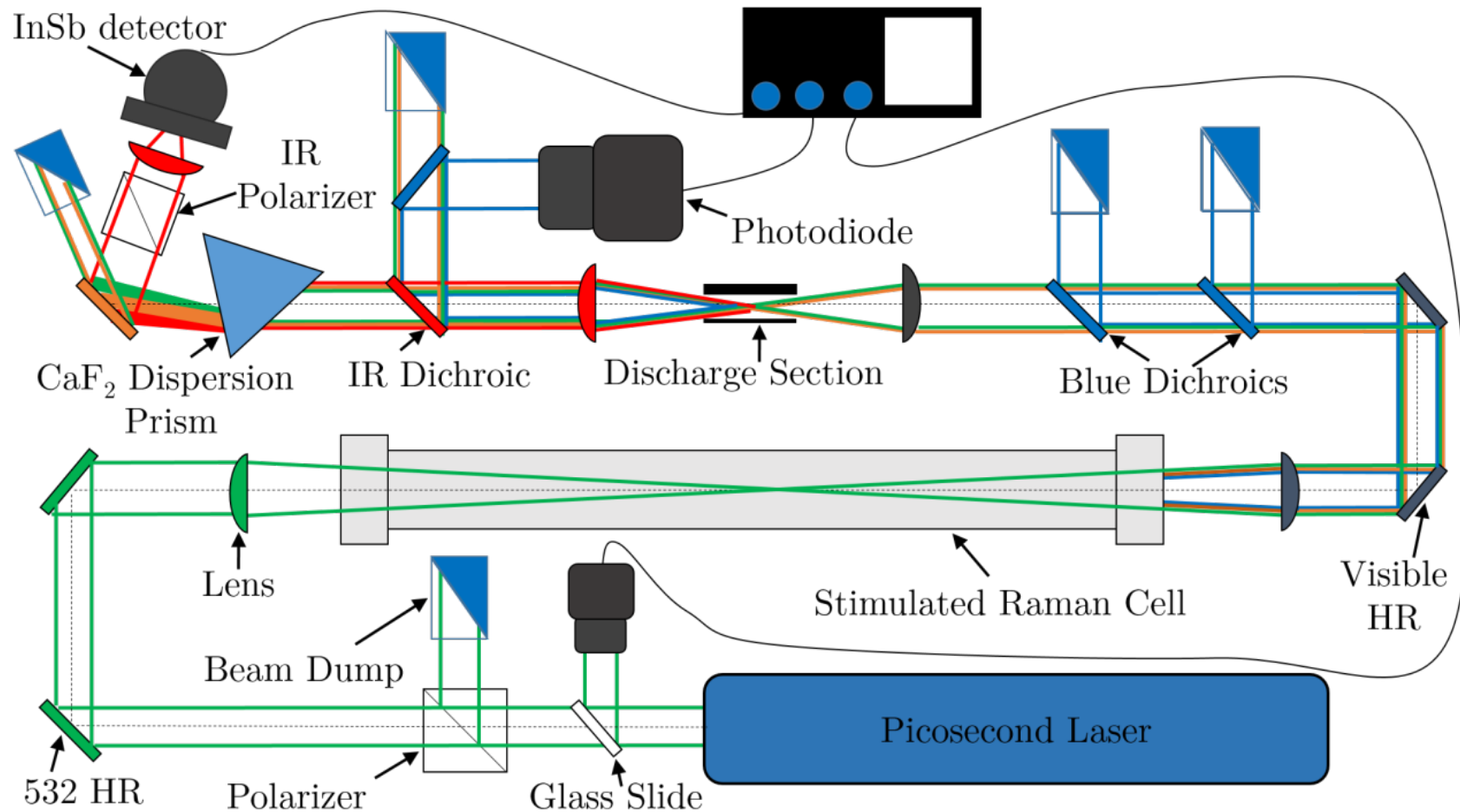


$$I_{\text{IR}} \sim \chi_{\text{IR}}^2 I_{\text{P}} I_{\text{Stokes}} E_{\text{Ext}}^2$$

$$|E_{\text{Ext}}| \sim \sqrt{\frac{\chi_{\text{CARS}}^2}{\chi_{\text{IR}}^2} \frac{I_{\text{IR}} I_{\text{P}}}{I_{\text{CARS}}}}$$

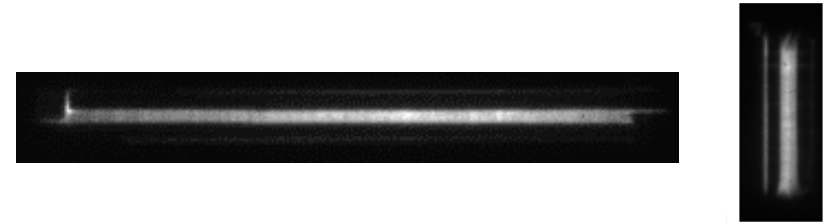
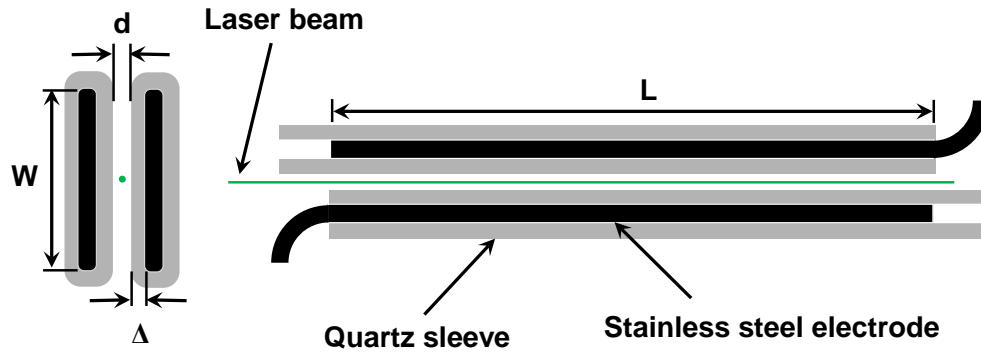
- Calibration using known electrostatic electric field
- Use of ps lasers: high peak laser intensity, higher 4-wave mixing signal

Ps 4-wave Mixing: Experimental Apparatus



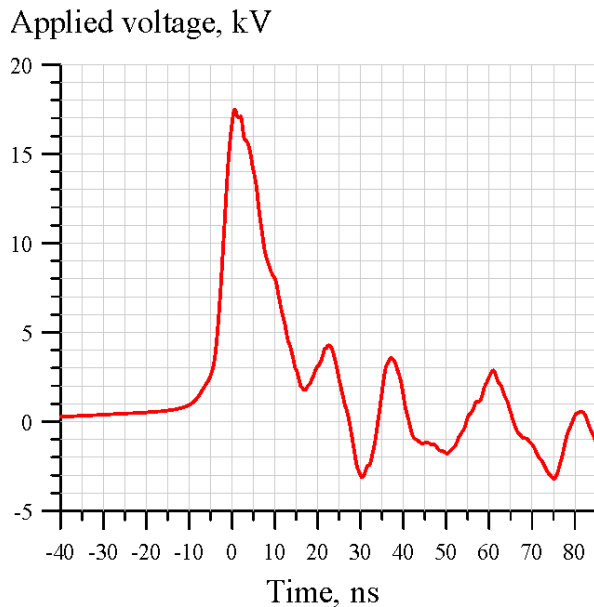
- **Key difference from previous work: use of stimulated Raman cell instead of a dye laser**
- **Raman cell filled with 15 bar N₂-He mixture**
- **Time resolution ~ 200 ps, spatial resolution (along laser beam) several cm**

Test Drive: Ns Pulse Discharge in Ambient Air Between Two Parallel Plate Electrodes



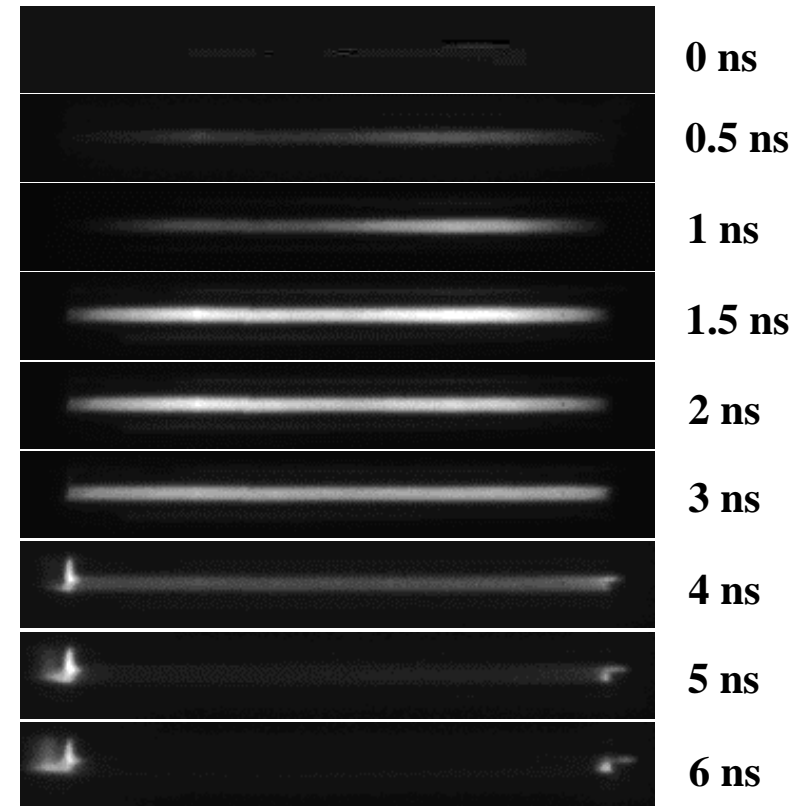
Camera gate 20 ns, single pulse

Two parallel plate electrodes encased in quartz sleeves



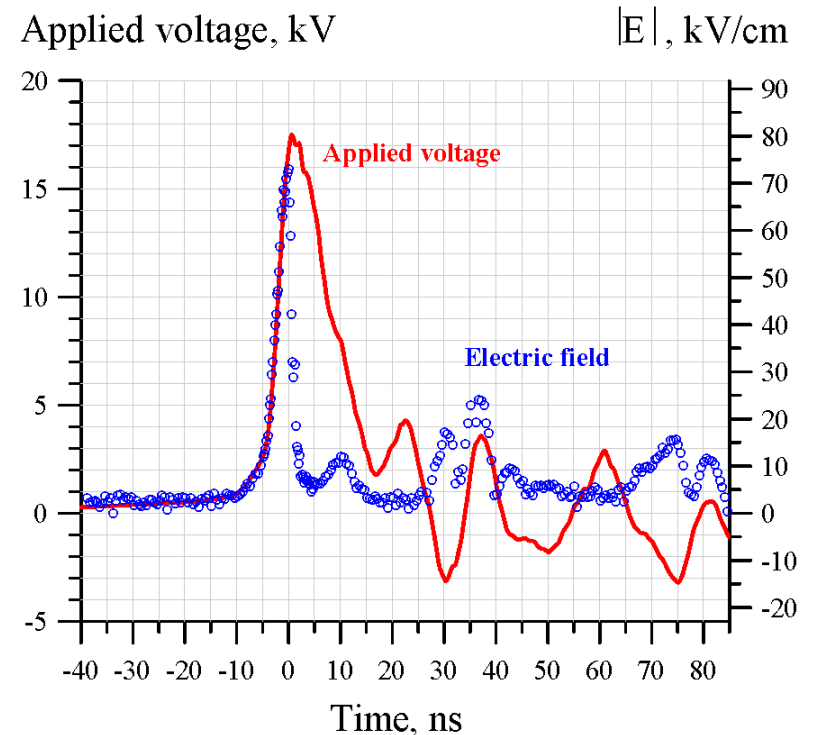
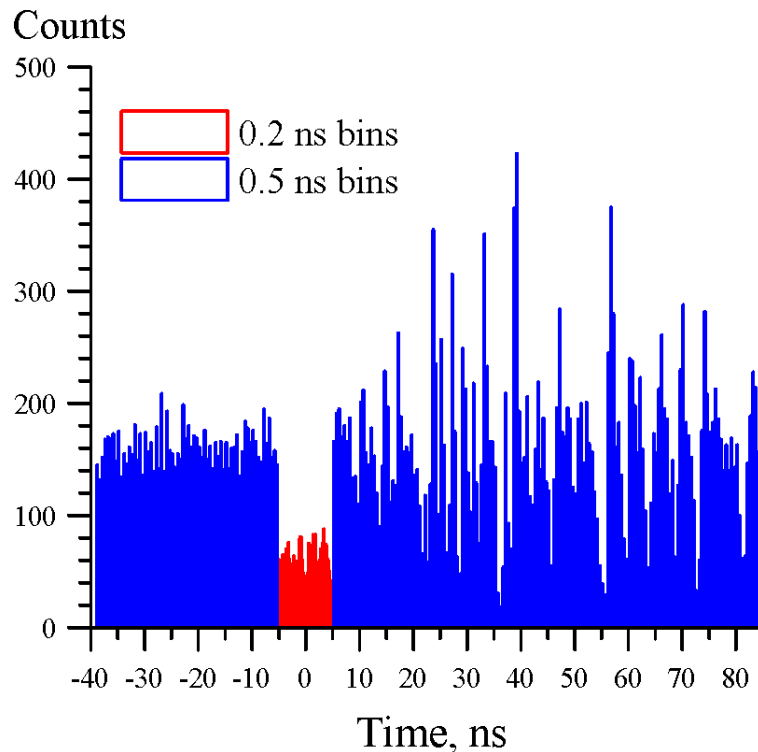
$dU/dt \approx 3 \text{ kV/ns}$, discharge gap 1.2 mm

Breakdown at $t \approx 0$ ns, $U \approx 17$ kV



Camera gate 1 ns, 50-pulse average

Time-Resolved Electric Field Measured with Sub-Ns Resolution (0.2-0.5 ns)



Laser shots (150 ps duration)
distributed randomly over 100 ns period

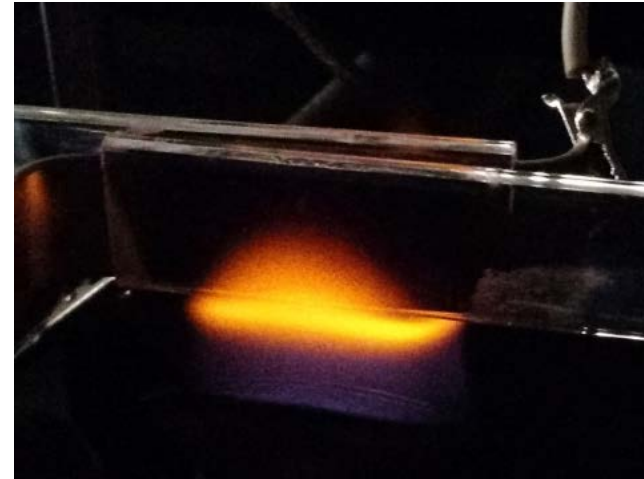
Four-wave mixing (electric field) data points
grouped into “time bins” 0.2-0.5 ns wide

- Electric field follows applied voltage until breakdown at 75 kV/cm, then falls rapidly
- Field after breakdown controlled by electron attachment, electron and ion transport

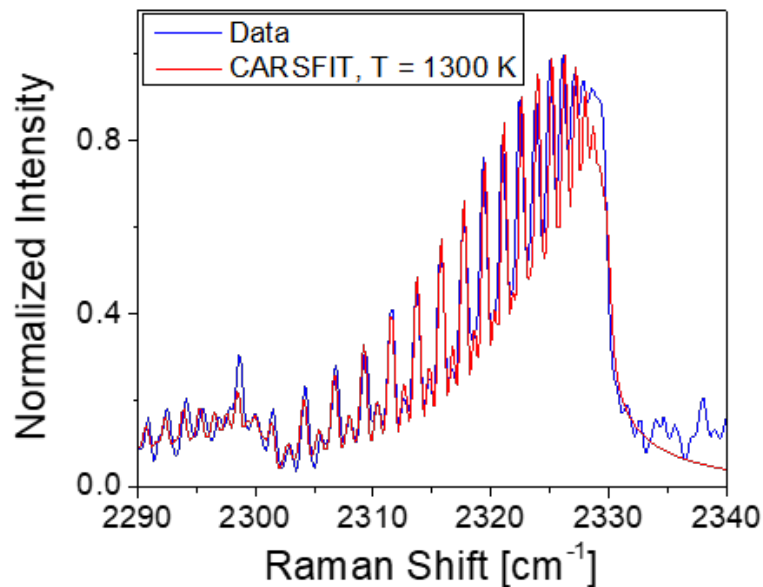
Hydrogen Diffusion Flame Between Two Parallel Plate Electrodes



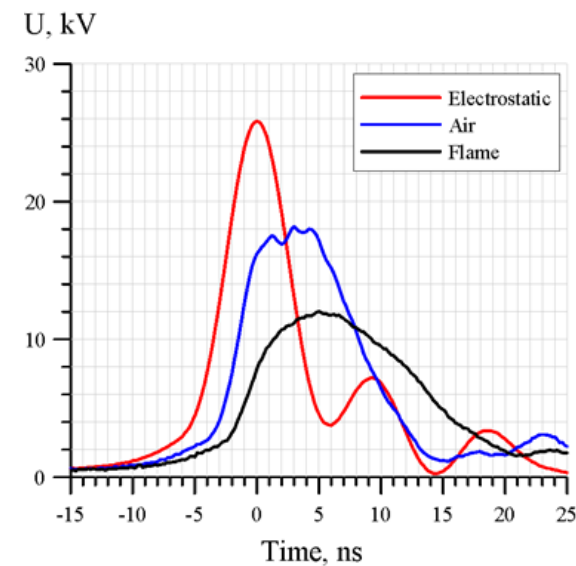
Hydrogen Diffusion Flame Setup



Hydrogen Flame

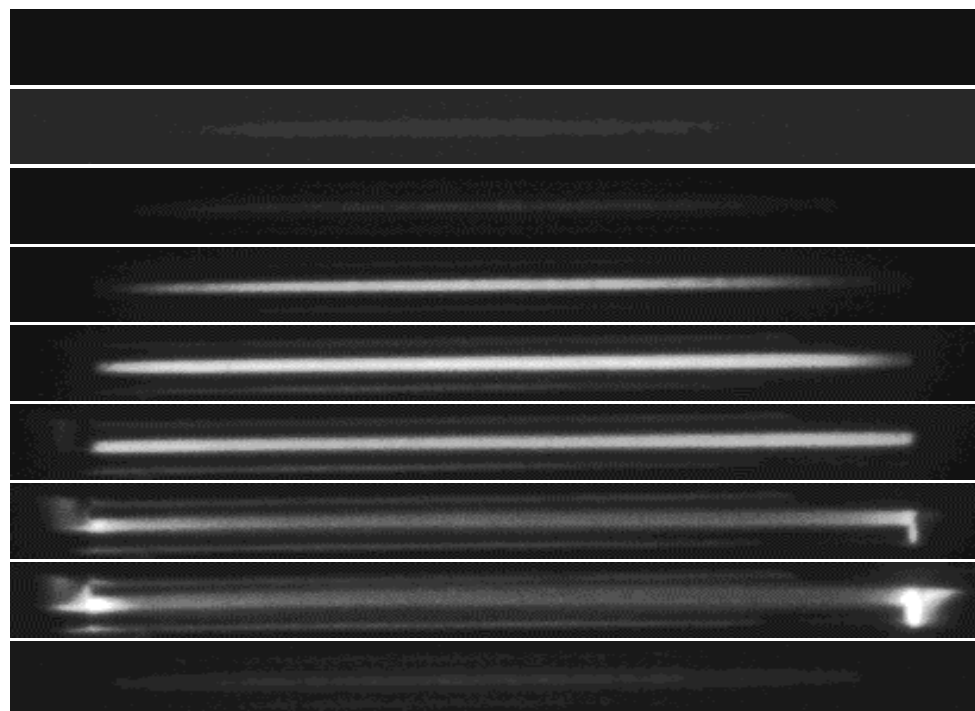


N_2 CARS spectrum: $T=1300\pm 50$ K

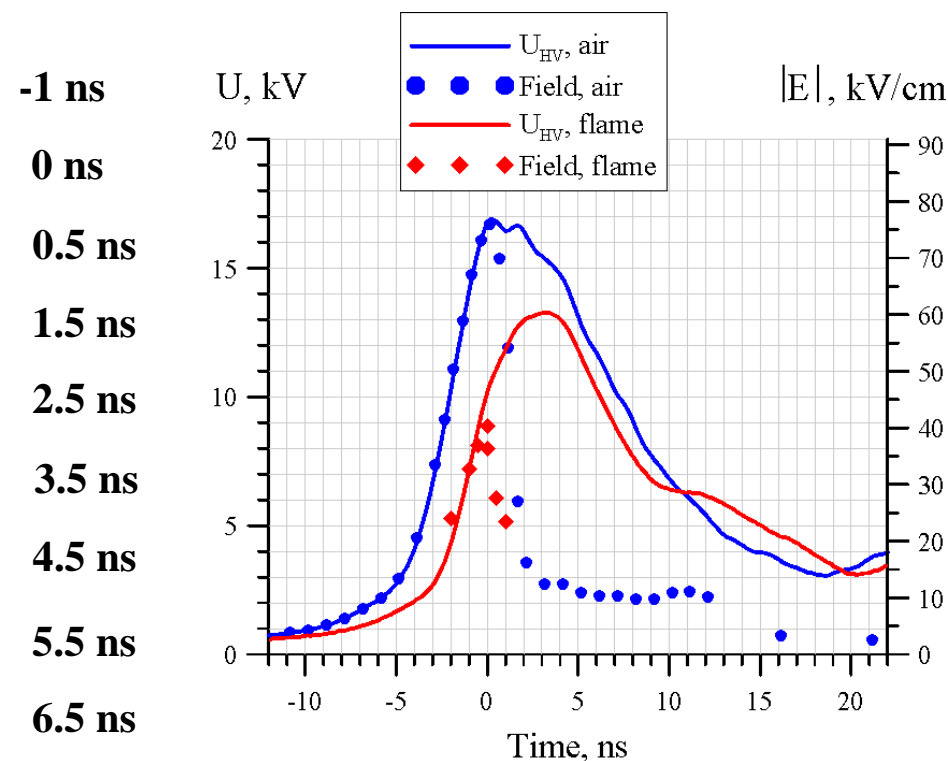


Voltage waveforms, discharge gap 1 mm

Comparison of Time-Resolved Electric Field in Ambient Air and in Hydrogen Flame



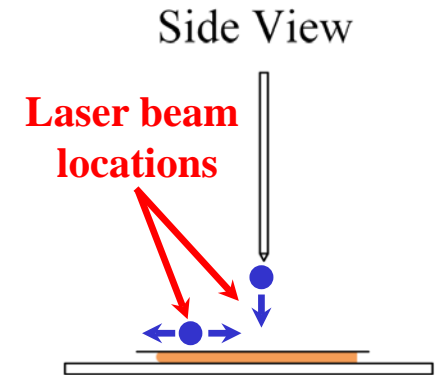
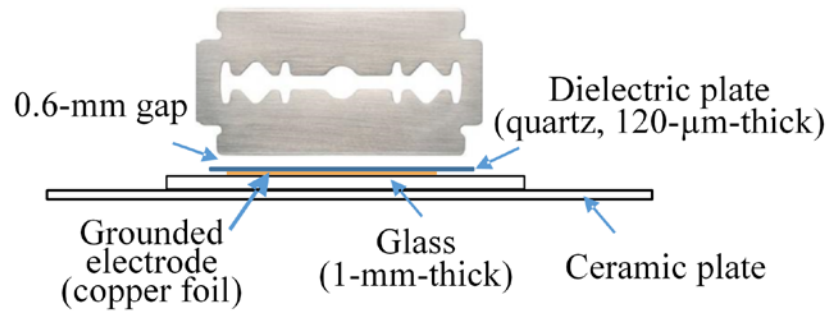
Top view, camera gate 1 ns, 50-pulse average



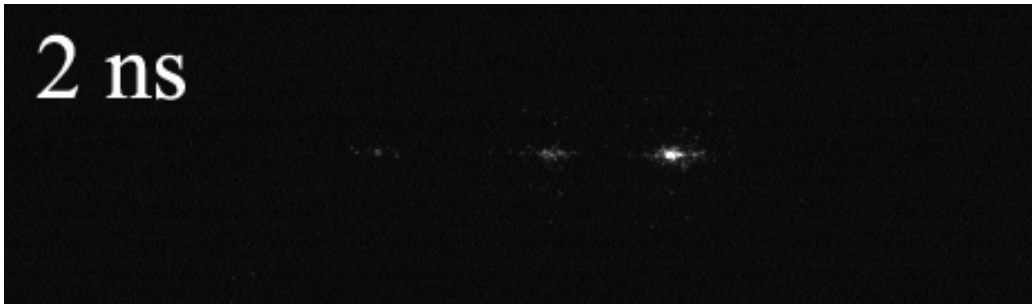
Discharge gap 1.3 mm

- Four-wave mixing signal in flame is 2 orders of magnitude weaker than in air
- Higher temperature, lower N_2 fraction in mixture, signal spread over many rotational states
- Breakdown field in flame a factor of ~ 2 lower than in room temperature air
- In both cases, field is higher than DC breakdown field, $E_{DC} \approx 30$ kV/cm

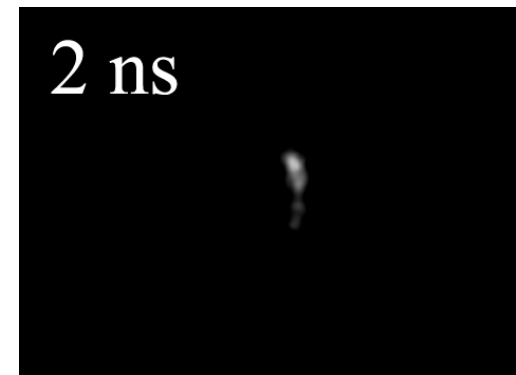
Negative Polarity Ns Pulse Discharge Over Quartz Surface



Top view, 2 ns gate

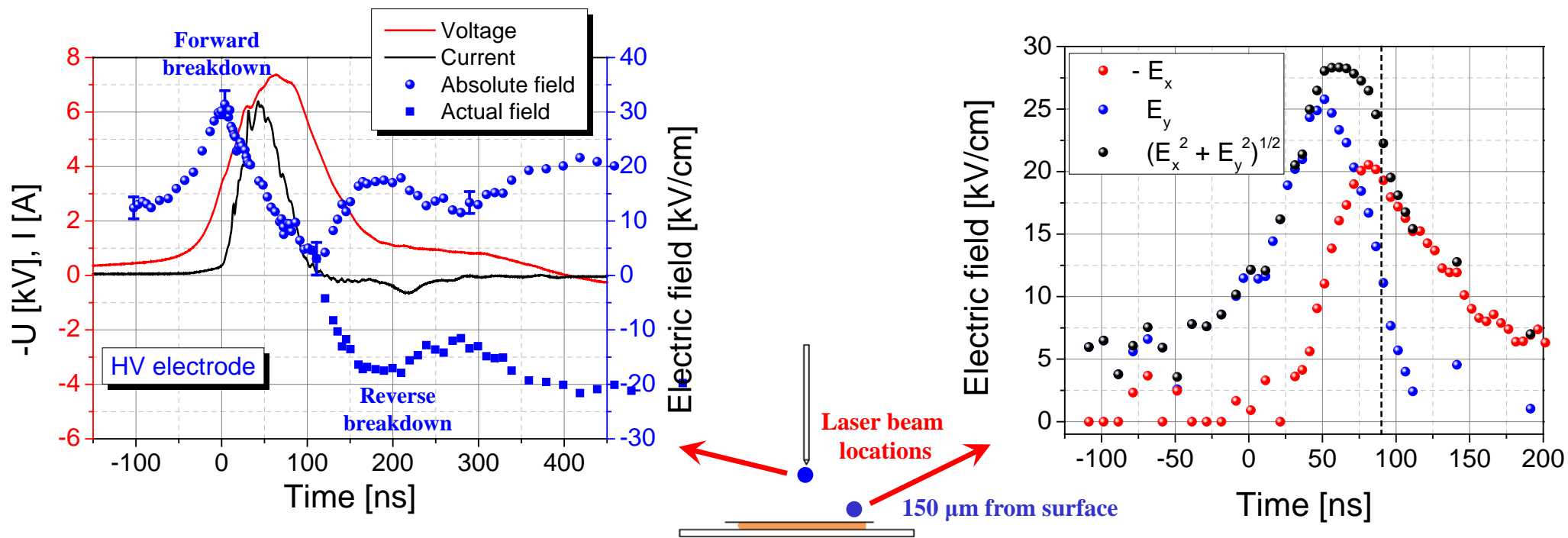


Side view, 2 ns gate



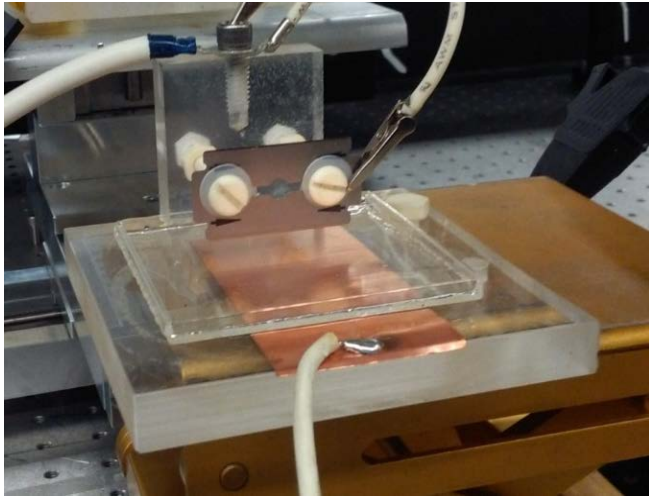
- Ns pulse discharge between a high-voltage electrode and a thin quartz plate
- Discharge gap 0.6 - 1.0 mm, diffuse surface ionization wave plasma ~ 200 μm thick
- Time-resolved electric field measured at multiple locations in the discharge gap

Electric Field Vector Components in a Surface Ionization Wave Discharge



- Initial field offset (at $t < 0$): charge accumulation on dielectric from previous pulse
- Breakdown field much lower than in “short” pulse discharge ($dU/dt = 0.1 \text{ kV/ns}$ vs. 3 kV/ns)
- After breakdown, field reduced due to charge accumulation on dielectric surface
- Field is reversed after applied voltage starts decreasing
- Away from HV electrode, field peaks later (E_y before E_x): surface ionization wave

Negative Polarity Ns Pulse Discharge Over Liquid Water

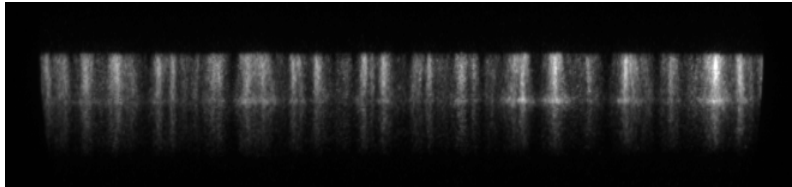
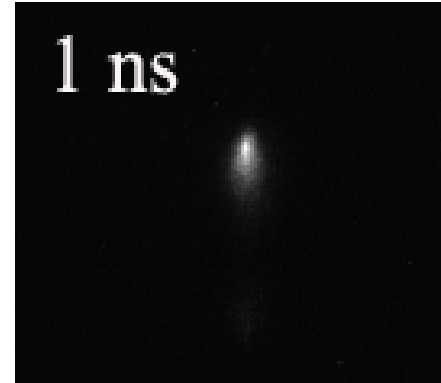


50-pulse accumulation movies

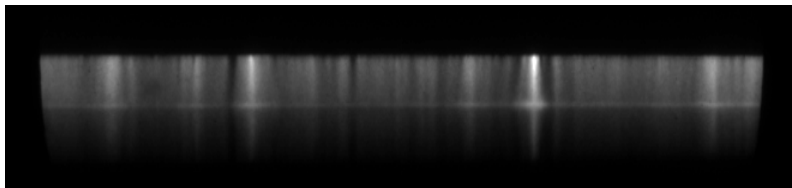
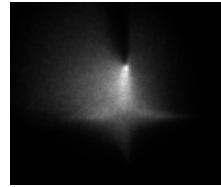
1 ns



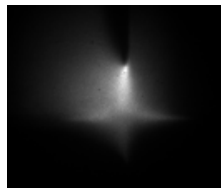
1 ns



Single-shot, 100 ns gate

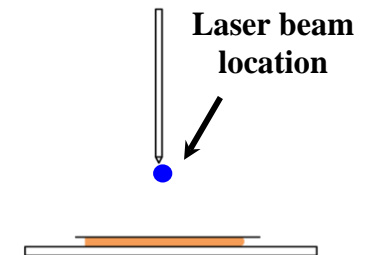
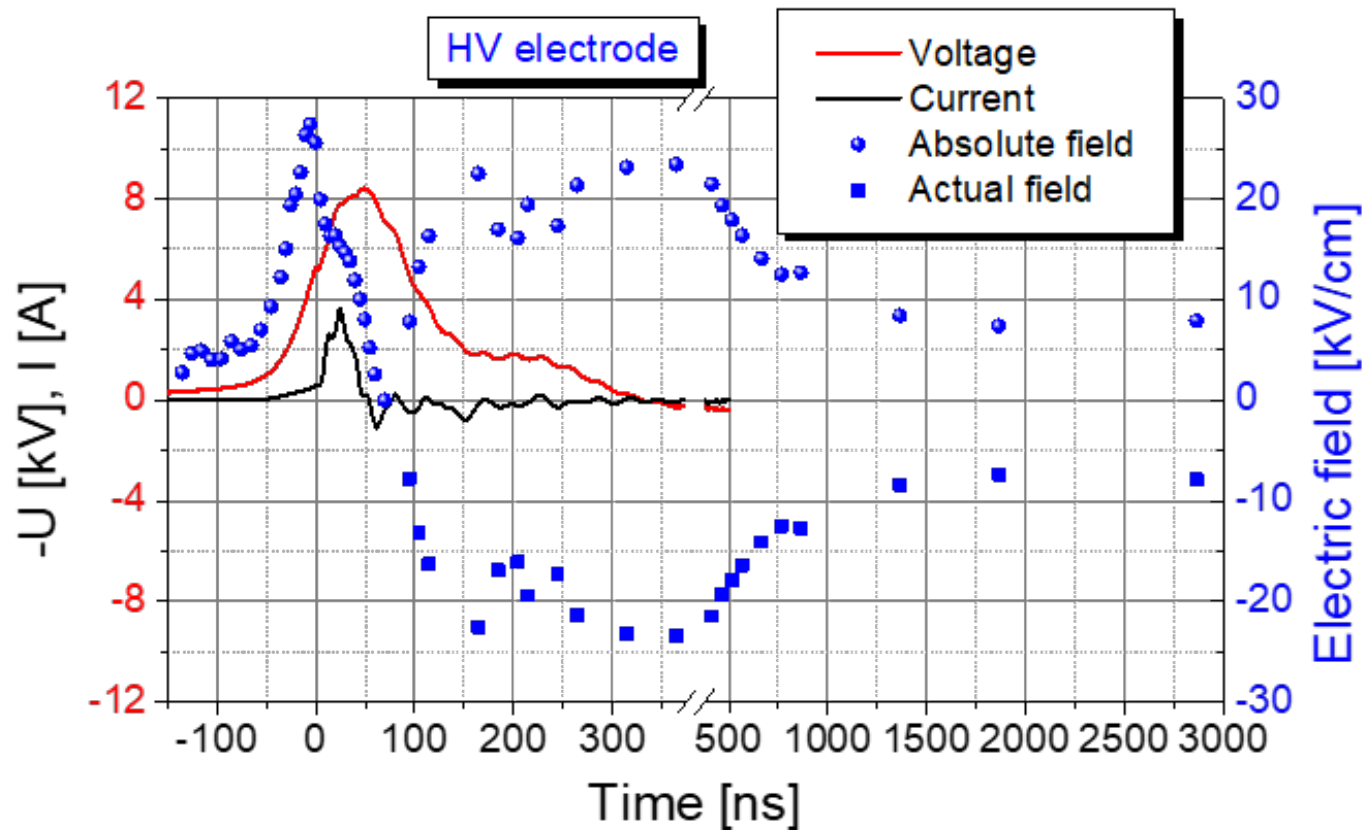


100-pulse accumulation



Positive polarity discharge plasma
is strongly filamentary

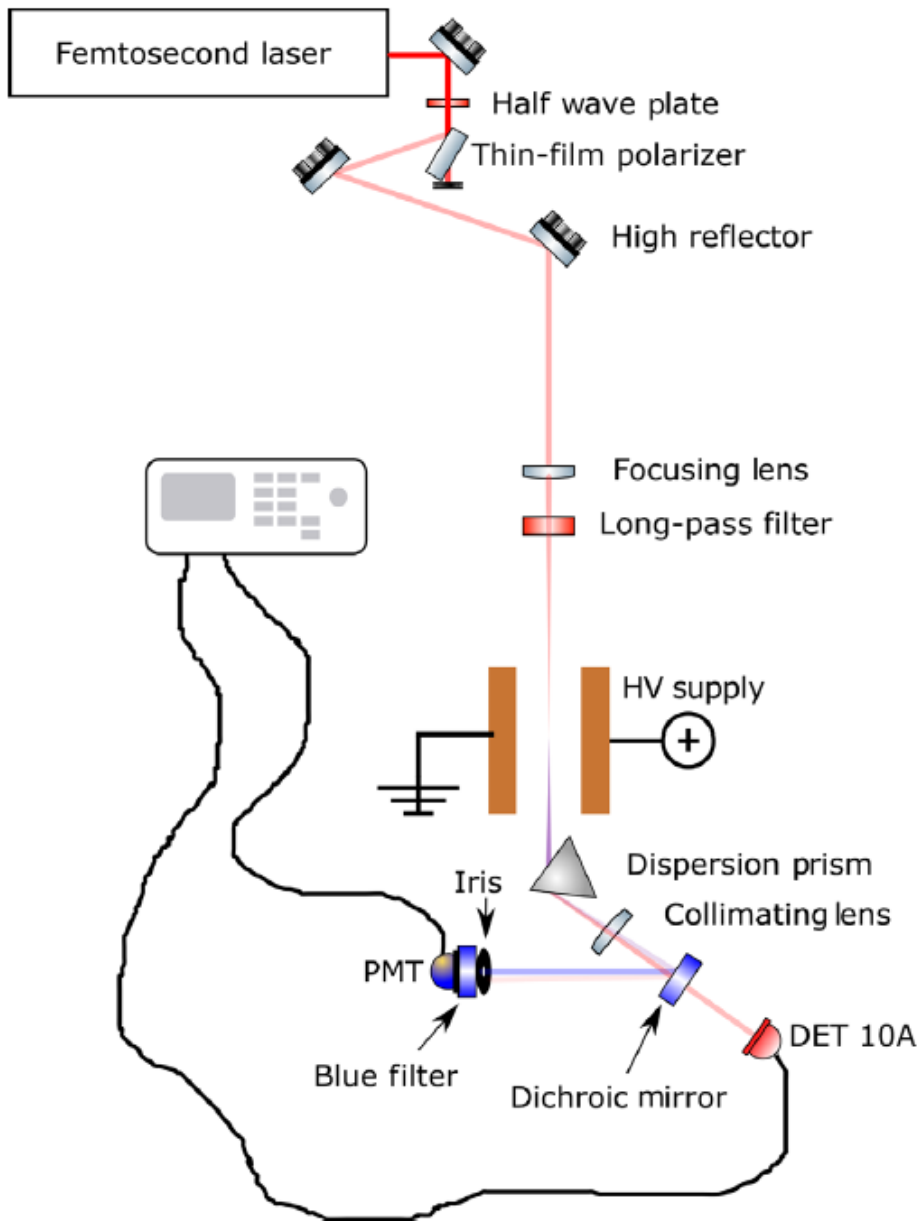
Electric Field in Discharge Over Liquid Water



- Results similar to measurements over quartz surface
- Field reduction to near zero over tens of ns after breakdown: plasma self-shielding
- Field reversal during voltage decay: surface charge accumulation
- Field decay on microsecond time scale: surface charge neutralization by ion transport

New Diagnostics for Electric Field Measurements: Femtosecond Second Harmonic Generation

Dogariu et al, Phys. Rev. Appl. 7 (2017) 024024



Polarization at 2ω induced by applied electric field, $E^{(F)}$, in presence of laser field, $E^{(\omega)}$:

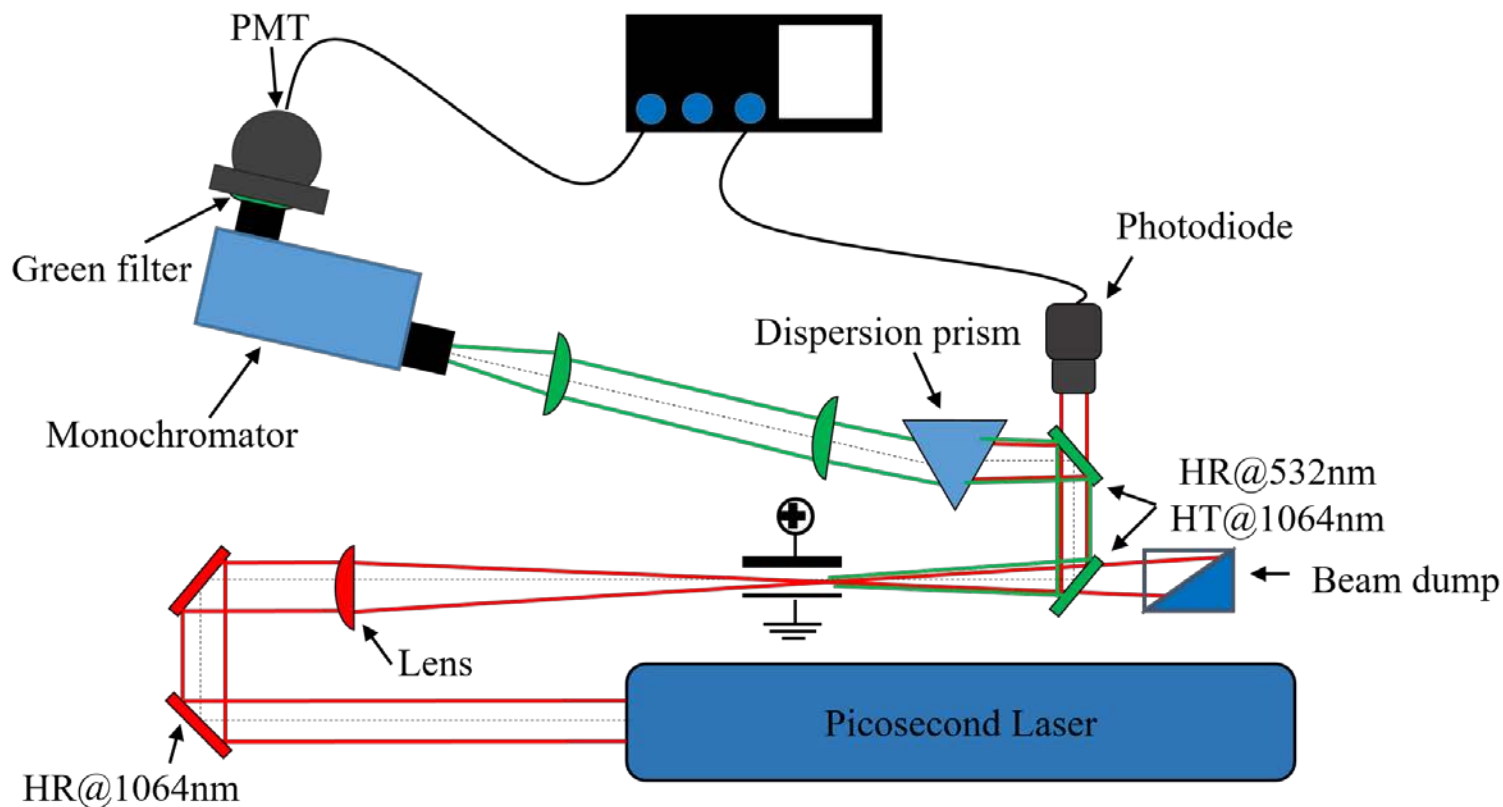
$$P_i^{(2\omega)} = \frac{3}{2} N \chi_{i,j,k,l}^{(3)}(-2\omega, 0, \omega, \omega) E_j^{(F)} E_k^{(\omega)} E_l^{(\omega)}$$

Signal proportional to laser intensity:
need an ultra-short-pulse (ps or fs) laser

Field parallel or perpendicular to laser polarization plane can be measured

Critical advantage: non-species sensitive diagnostics, would work in reacting mixtures and flames

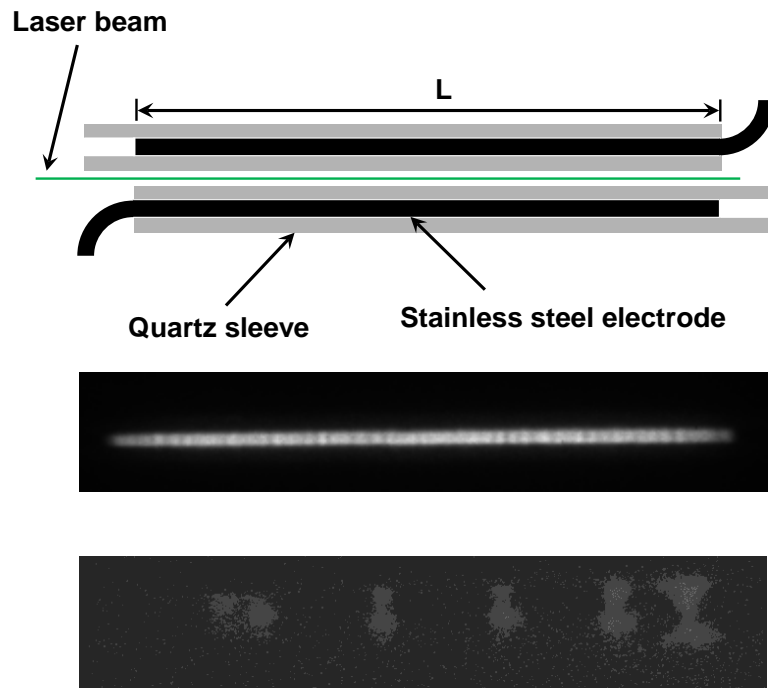
Picosecond Second Harmonic Generation Experiment



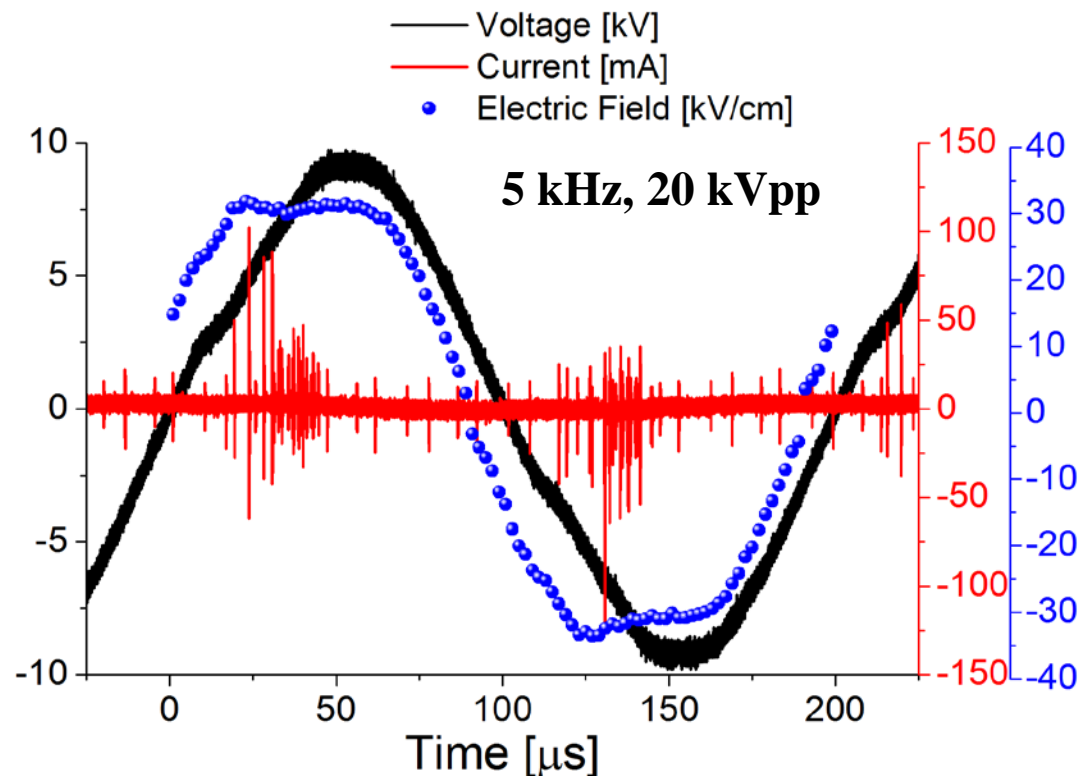
- Laser pulse 30 ps long, 2-4 mJ/pulse at 1064 nm, generating second harmonic at 532 nm
- **SHG vs. compared to four-wave mixing: 10 times more signal at 10 times less laser power**
- Goal: electric field measurements in fuel-air plasmas and atmospheric pressure plasma jets
- Test cases: AC Dielectric Barrier Discharge (DBD) plasmas

Preliminary Results:

Electric Field in Plane-to-Plane AC DBD plasma



Microdischarges (not to scale)

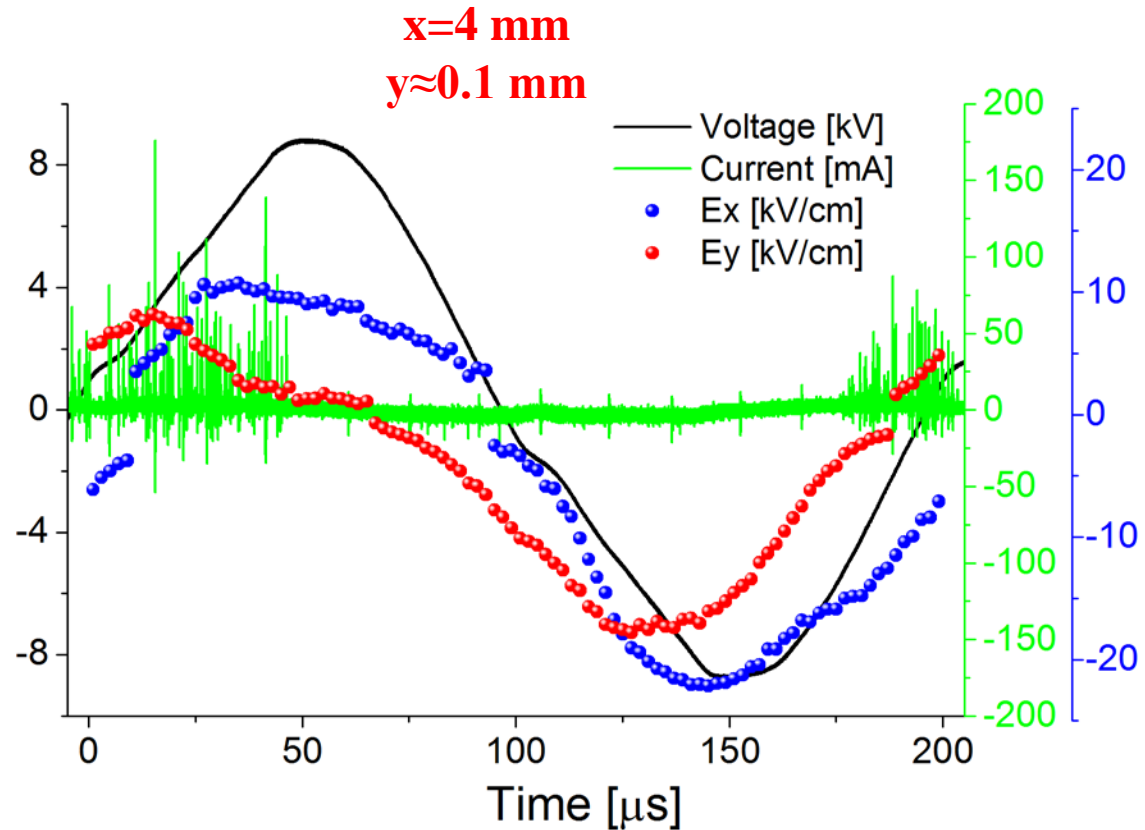
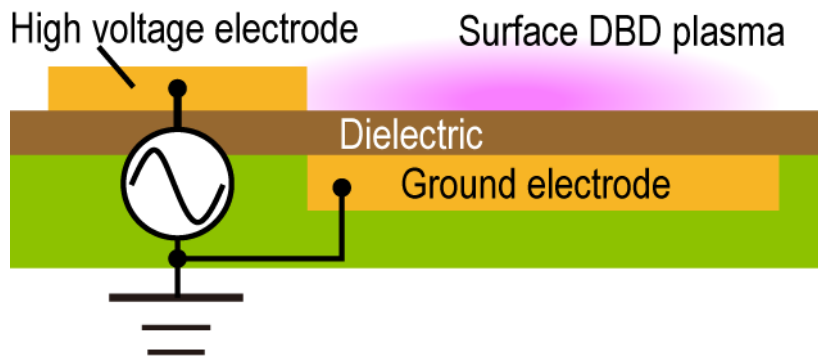
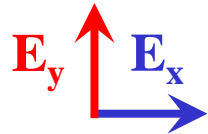
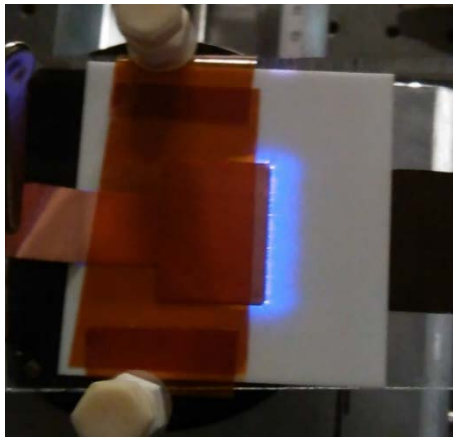


Test case I: AC DBD plasma in room air, two parallel plane electrodes 50 mm long, 5 kHz frequency, 20 kV peak-to-peak voltage, two dielectric barriers (quartz sleeves 1.5 mm thick), discharge gap 1.3 mm

Phase shift between applied voltage and electric field, due to surface charges

Micro-discharges keep electric field near breakdown threshold (≈ 30 kV/cm)

Preliminary Results: Electric Field in AC Surface DBD Plasma Actuator



Test case II: AC DBD surface plasma actuator in room air, alumina dielectric plate 1.5 mm thick, electrodes 25 mm wide, 5 kHz frequency, 17 kV peak-to-peak voltage

Electric field vector components (E_x and E_y) measured at $x=0-9$ mm from high voltage electrode, $y \approx 100 \mu\text{m}$ above dielectric surface

Quantitative insight into mechanism of plasma actuator interaction with the flow?

Summary

- Electric field vector is measured by ps 4-wave mixing in ns pulse discharges in ambient air and in H₂-air flame
- Demonstrated capability for sub-ns time resolution, measurements of individual electric field vector components
- Breakdown field in air and in flame considerably exceeds DC breakdown threshold, for pulse voltage rise time of ~ 5 ns
- Time-resolved, spatially resolved electric field (E_x and E_y) measured in ns pulse, surface ionization wave plasmas in ambient air, over quartz and water surfaces
- Ps second harmonic generation diagnostics for electric field measurements has been tested in AC DBD plasmas; further measurements underway